Results in Mathematics



Non-coercive Neumann Boundary Control Problems

Thomas Apel, Mariano Mateos, and Arnd Rösch

Abstract. The article examines a linear-quadratic Neumann control problem that is governed by a non-coercive elliptic equation. Due to the nonself-adjoint nature of the linear control-to-state operator, it is necessary to independently study both the state and adjoint state equations. The article establishes the existence and uniqueness of solutions for both equations, with minimal assumptions made about the problem's data. Next, the regularity of these solutions is studied in three frameworks: Hilbert-Sobolev spaces, Sobolev-Slobodeckiĭ spaces, and weighted Sobolev spaces. These regularity results enable a numerical analysis of the finite element approximation of both the state and adjoint state equations. The results cover both convex and non-convex domains and quasi-uniform and graded meshes. Finally, the optimal control problem is analyzed and discretized. Existence and uniqueness of the solution, first-order optimality conditions, and error estimates for the finite element approximation of the control are obtained. Numerical experiments confirming these results are included.

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Thomas Apel, Mariano Mateos and Arnd Rösch have contributed equally to this work.

1. Introduction

Let us consider a domain $\Omega \subset \mathbb{R}^2$ with a polygonal boundary Γ . We are concerned with the Neumann boundary control problem

(P)
$$\min_{u \in U_{ad}} J(u) := \frac{1}{2} \int_{\Omega} (y_u(x) - y_d(x))^2 \, \mathrm{d}x + \frac{\nu}{2} \int_{\Gamma} u^2(x) \, \mathrm{d}x + \int_{\Gamma} y_u(x) g_{\varphi}(x) \, \mathrm{d}x$$

where $y_d \in L^2(\Omega)$ and $g_{\varphi} \in L^2(\Gamma)$ are given functions, $\nu > 0$,

$$U_{\rm ad} = \{ u \in L^2(\Omega) : u_a \le u(x) \le u_b \text{for a.e.} x \in \Omega \}$$

with $-\infty \leq u_a < u_b \leq +\infty$, and y_u is the solution of

$$\begin{cases} Ay + b(x) \cdot \nabla y + a_0(x)y = 0 & \text{in } \Omega, \\ \partial_{n_A} y = u & \text{on } \Gamma. \end{cases}$$
(1.1)

Assumptions regarding the symmetric second order differential operator A and the coefficients b and a_0 will be described later. Let us just emphasize now that we will make no assumptions on b and a_0 that would imply coerciveness of the associated bilinear form.

The main objective of this paper is to discretize the optimal control problem using the finite element method and to obtain error estimates for the approximations of the optimal control in terms of the discretization parameter h. The paper aims to minimize assumptions to better capture their essence. The results are valid for possibly non-convex domains and both quasi-uniform and graded meshes. Although the theory for Neumann boundary optimal control problems governed by elliptic equations is quite complete, to the best of our knowledge, the issues that arise when the elliptic operator governing the equation is not coercive have not been addressed yet; see Casas, Mateos and Tröltzsch 2005 [1], Casas and Mateos 2007 [2], Mateos and Rösch 2011 [3], Apel Pfefferer and Rösch 2012 and 2015 [4,5], Krumbiegel and Pfefferer 2015 [6] or the thesis by Winkler 2015 [7]. The only papers, we are aware, that deal with optimal control problems governed by a non-coercive elliptic equation are about distributed controls; see Casas, Mateos and Rösch 2020 and 2021 [8,9]. In both papers, this fact and the convexity of the domain are used in an essential way in some of the proofs, and hence those results are not applicable to our problem.

We will see that problem (P) has a unique solution \bar{u} , and that it satisfies the optimality conditions, which we state now in an informal way: there exist \bar{y} and $\bar{\varphi}$ such that

$$\begin{cases} A\bar{y} + b(x) \cdot \nabla \bar{y} + a_0(x)\bar{y} = 0 & \text{ in } \Omega, \\ \partial_{n_A}\bar{y} = \bar{u} & \text{ on } \Gamma, \end{cases}$$
(1.2a)

$$\begin{cases} A\bar{\varphi} - \nabla \cdot (b(x)\bar{\varphi}) + a_0(x)\bar{\varphi} = \bar{y} - y_d & \text{in } \Omega, \\ \partial_{n_A}\bar{\varphi} + \varphi b \cdot n = g_{\varphi} & \text{on } \Gamma, \end{cases}$$
(1.2b)

$$\int_{\Gamma} (\bar{\varphi} + \nu \bar{u})(u - \bar{u}) \, \mathrm{d}x \ge 0 \,\,\forall u \in U_{\mathrm{ad}}.$$
(1.2c)

Since (P) is a linear-quadratic strictly convex problem, existence and uniqueness of the solution follow in a standard way once we have proved existence and uniqueness of solution of the state equation and continuity of the controlto-state mapping. But, since we will not formulate any assumptions on b or a_0 that would lead to a coercive operator, this task is not standard. In particular, div b may be large, such that the usual assumption $a_0 - \frac{1}{2} \text{div } b \ge c_0 > 0$ is not satisfied. This will be done in Sect. 2.

In Sect. 3 we investigate the regularity properties of the solutions of the state equation and the adjoint state equations. Since these are different, we perform this task in two steps resulting in Theorems 3.4 and 3.5, respectively. We obtain results in Hilbert–Sobolev, in Sobolev–Slobodeckiĭ and in weighted Sobolev spaces, with our focus on treating the numerical approximation of (P) in non-convex domains. The regularity results in non-weighted spaces serve us as intermediate results to prove the error estimates in weighted Sobolev spaces, but they are also of independent interest. Note that, although regularity results for elliptic boundary value problems are widely investigated, see, e. g., the monographs [10–14], the particular results which we need for our approximation theory were not available for non-coercive problems with variable coefficients.

In Sect. 4 we study the numerical discretization of both the state and adjoint state equation. We obtain existence and uniqueness of the solution as well as error estimates. Our results are valid in convex and non-convex domains and for quasi-uniform and graded meshes, with possibly a non-optimal grading parameter μ .

With these results at hand, we will be able to deduce existence, uniqueness, and optimality conditions in Sect. 5. Moreover, regularity properties of the optimal solution and its related state and adjoint state are given in terms of weighted Sobolev spaces. Finally, we will discretize the control problem. The control is approximated using piecewise constant functions whereas the state and adjoint state are discretized by continuous piecewise linear functions. A close inspection of the proofs in the above mentioned papers about Neumann control problems, suggests that, if no postprocessing step is done, the order of convergence of the error in $L^2(\Gamma)$ for the control variable will be limited by the order of convergence of the finite element error in $H^1(\Omega)$ for the state or the adjoint state equation; see e.g. the proof of Lemma 4.7 in [1]. This means that, for a non-convex domain and a quasi-uniform mesh, the order of convergence that can be obtained—applying the usual techniques in optimal control together with the regularity results and the finite element error estimates provided in this paper—is approximately h^{λ} , where $1/2 < \lambda < 1$. For instance, in the problem shown as an example in Sect. 6, $h^{2/3}$ would be expected. Nevertheless, the numerical experiments show clearly order h, and we are able to get that in Theorem 5.7: If the corner singularities are of type r^{λ_j} , the index j counting the corners, and the mesh is graded near the corners with parameter μ_j , then the approximation order of the control is $s^* \leq 1$ with $s^* < \frac{3\lambda_j}{2\mu_j}$, i. e., $s^* = 1$ is achieved if $\mu_j < \frac{3}{2}\lambda_j$ for all j. In the works by Apel, Pfefferer and Rösch [4,5] a stronger grading $\mu_j < \lambda_j$ is used to obtain an optimal control of convergence for the so-called post-processed control, i.e., the pointwise projection onto the admissible set of $-\bar{\varphi}_h/\nu$, where $\bar{\varphi}_h$ is the discrete adjoint state associated to the discrete optimal control. In Theorem 4.2.1 of the thesis of Winkler [7] it is shown that for quasi-uniform meshes, i.e., $\mu_j = 1$, the order $s^* = 1$ is achieved for any angle.

2. Existence, Uniqueness and Continuous Dependence of the Solution of the State and Adjoint State Equations

On A, b and a_0 we make the following assumptions.

Assumption 2.1. A is the operator given by

$$Ay = -\sum_{i,k=1}^{2} \partial_{x_k} (a_{ik}(x)\partial_{x_i}y) \text{ with } a_{ik} \in L^{\infty}(\Omega),$$

 $a_{ik} = a_{ki}$ for $1 \le i, k \le 2$, and satisfying the following ellipticity condition:

$$\exists \Lambda > 0 \text{ such that } \sum_{i,k=1}^{2} a_{ik}(x)\xi_i\xi_k \ge \Lambda |\xi|^2 \quad \forall \xi \in \mathbb{R}^2 \text{ and for a.a. } x \in \Omega.$$
 (2.1)

The function $b: \Omega \to \mathbb{R}^2$ satisfies $b \in L^{\hat{p}}(\Omega)^2$ with $\hat{p} > 2$ and there exists $\hat{q} > 1$ such that $\nabla \cdot b \in L^{\hat{q}}(\Omega)$ and $b \cdot n \in L^{\hat{q}}(\Gamma)$. For the function $a_0: \Omega \to \mathbb{R}$ it is assumed that $a_0 \in L^{\hat{q}}(\Omega), a_0(x) \ge 0$ for a.e. $x \in \Omega$ and there exists $E \subset \Omega$ with |E| > 0 such that $a_0(x) \ge \Lambda/2$ for all $x \in E$.

Remark 2.2. Note that this assumption does not lead to a coercive bilinear form. While Assumption 2.1 is sufficient for the purposes of proving existence and uniqueness of solution, to establish adequate regularity results for the solution, further regularity must be imposed on the coefficients. The reader is referred to the results of Sect. 3 for the details required in the different scenarios.

Before addressing the main results of this section, we recall some well known inequalities that will be used throughout this paper.

We will often use the following form of Hölder's inequality: for $q, p_1, \ldots, p_k \in [1, \infty]$ such that $1/p_1 + \cdots + 1/p_k \leq 1/q$ and $f_i \in L^{p_i}(\Omega), i = 1, \ldots, k$ there exists a constant $C_{\Omega} = |\Omega|^{1/q - (1/p_1 + \cdots + 1/p_k)}$, such that $||f_1 \cdots f_k||_{L^q(\Omega)} \leq C_{\Omega} ||f_1||_{L^{p_1}(\Omega)} \cdots ||f_k||_{L^{p_k}(\Omega)}$.

The inequality

$$\|y\|_{H^1(\Omega)} \le C_E(\|\nabla y\|_{L^2(\Omega)} + \|y\|_{L^2(E)}) \quad \forall y \in H^1(\Omega)$$
(2.2)

is a generalization of Poincaré's inequality and can be found, e.g., in [15, Theorem 11.19]. In dimension 2, Sobolev's embedding theorem gives that for all $r < \infty$ there exists $K_{\Omega,r} > 0$ such that

$$\|y\|_{L^{r}(\Omega)} \leq K_{\Omega,r} \|y\|_{H^{1}(\Omega)} \quad \forall y \in H^{1}(\Omega).$$

$$(2.3)$$

We will denote by $\langle \cdot, \cdot \rangle_{\Omega}$ the duality product in $H^1(\Omega)' \times H^1(\Omega)$ and by $\langle \cdot, \cdot \rangle_{\Gamma}$ the duality product in $H^{1/2}(\Gamma)' \times H^{1/2}(\Gamma)$. We notice that any $g \in$ $H^{1/2}(\Gamma)'$ defines an element in $H^1(\Omega)'$, which will be denoted in the same way by

$$\langle g, z \rangle_{\Omega} = \langle g, \operatorname{tr} z \rangle_{\Gamma} \ \forall z \in H^1(\Omega).$$
 (2.4)

In this case, we will simply write $\langle g, z \rangle_{\Gamma}$. Also notice that for any fixed q > 1, the functions $f \in L^q(\Omega)$ and $g \in L^q(\Gamma)$ define elements in $H^1(\Omega)'$ and $H^{1/2}(\Gamma)'$ respectively by

$$\langle f, z \rangle_{\Omega} = \int_{\Omega} f z \, \mathrm{d}x, \qquad \langle g, z \rangle_{\Gamma} = \int_{\Gamma} g z \, \mathrm{d}x \,\,\forall z \in H^1(\Omega).$$
 (2.5)

For every $y \in H^1(\Omega)$, we define $\mathcal{A}y$ by

$$\langle \mathcal{A}y, z \rangle_{\Omega} = \int_{\Omega} \sum_{i,k=1}^{2} a_{ik} \partial_{x_{i}} y \partial_{x_{k}} z \, \mathrm{d}x + \int_{\Omega} (b \cdot \nabla y) z \, \mathrm{d}x + \int_{\Omega} a_{0} y z \, \mathrm{d}x \, \forall z \in H^{1}(\Omega).$$

$$(2.6)$$

Using this operator, we have that the weak form of the state Eq. (1.1) is: find $y_u \in H^1(\Omega)$ such that

$$\langle \mathcal{A}y_u, z \rangle_{\Omega} = \langle u, z \rangle_{\Gamma} \ \forall z \in H^1(\Omega).$$
 (2.7)

We first prove continuity of the operator \mathcal{A} and Gårding's inequality. We adapt the proof of [8, Lemma 2.1]

Lemma 2.3. Under Assumption 2.1 we have that $\mathcal{A} \in \mathcal{L}(H^1(\Omega), H^1(\Omega)')$ and there exists a constant $C_{\Lambda, E, b}$ such that

$$\langle \mathcal{A}z, z \rangle_{\Omega} \ge \frac{\Lambda}{8C_E^2} \|z\|_{H^1(\Omega)}^2 - C_{\Lambda, E, b} \|z\|_{L^2(\Omega)}^2 \quad \forall z \in H^1(\Omega)$$
(2.8)

where Λ and C_E are the constants from (2.1) and (2.2), respectively.

Proof. Let us show that \mathcal{A} is a linear continuous operator. Denote $S = \{z \in H^1(\Omega) : \|z\|_{H^1(\Omega)} = 1\}$. We split \mathcal{A} into three parts $\mathcal{A}_i, i = 1, 2, 3$.

$$\begin{split} \|\mathcal{A}_{1}y\|_{H^{1}(\Omega)'} &= \sup_{z \in S} \int_{\Omega} \sum_{i,k=1}^{2} a_{ik} \partial_{x_{i}} y \partial_{x_{k}} z \, \mathrm{d}x \\ &\leq \sup_{z \in S} 4 \max_{1 \leq i,k \leq 2} \|a_{ik}\|_{L^{\infty}(\Omega)} \|\nabla y\|_{L^{2}(\Omega)^{2}} \|\nabla z\|_{L^{2}(\Omega)^{2}} \\ &\leq 4 \max_{1 \leq i,k \leq 2} \|a_{ik}\|_{L^{\infty}(\Omega)} \|y\|_{H^{1}(\Omega)}. \end{split}$$

Take now $s_p > 1$ such that $1/s_p = 1/\hat{p} + 1/2$. From (2.3) and Hölder's inequality we infer for every $y \in H^1(\Omega)$

$$\begin{aligned} \|\mathcal{A}_{2}y\|_{H^{1}(\Omega)'} &= \sup_{z \in S} \int_{\Omega} (b \cdot \nabla y) z \, \mathrm{d}x \le \|b \cdot \nabla y\|_{L^{s_{p}}(\Omega)} \|z\|_{L^{s'_{p}}(\Omega)} \\ &\le K_{\Omega,s'_{p}} \|b\|_{L^{\hat{p}}(\Omega)^{2}} \|\nabla y\|_{L^{2}(\Omega)^{2}} \le K_{\Omega,s'_{p}} \|b\|_{L^{\hat{p}}(\Omega)^{2}} \|y\|_{H^{1}(\Omega)}, \end{aligned}$$

Fix now some $s_q \in (1, \hat{q})$ and take $r \in (1, +\infty)$ such that $1/\hat{q} + 1/r = 1/s_q$. From (2.3) we infer that

$$\begin{aligned} \|\mathcal{A}_{3}y\|_{H^{1}(\Omega)'} &= \sup_{z \in S} \int_{\Omega} a_{0}yz \, \mathrm{d}x \le \|a_{0}y\|_{L^{s_{q}}(\Omega)} \|z\|_{L^{s'_{q}}(\Omega)} \\ &\le K_{\Omega,s'_{q}} \|a_{0}\|_{L^{\hat{q}}(\Omega)} \|y\|_{L^{r}(\Omega)} \le K_{\Omega,s'_{q}} K_{\Omega,r} \|a_{0}\|_{L^{\hat{q}}(\Omega)} \|y\|_{H^{1}(\Omega)}. \end{aligned}$$

Hence, we have that \mathcal{A} is a well-posed linear and continuous operator.

Let us prove (2.8). Using Assumption 2.1, (2.2), and Young and Hölder inequalities we get

$$\begin{split} \langle \mathcal{A}z, z \rangle_{\Omega} &\geq \Lambda \|\nabla z\|_{L^{2}(\Omega)^{2}}^{2} + \frac{\Lambda}{2} \|z\|_{L^{2}(E)}^{2} - \|\nabla z\|_{L^{2}(\Omega)^{2}} \|bz\|_{L^{2}(\Omega)^{2}} \\ &\geq \frac{\Lambda}{2} \|\nabla z\|_{L^{2}(\Omega)^{2}}^{2} + \frac{\Lambda}{2} \|z\|_{L^{2}(E)}^{2} - \frac{1}{2\Lambda} \|bz\|_{L^{2}(\Omega)^{2}}^{2} \\ &\geq \frac{\Lambda}{4C_{E}^{2}} \|z\|_{H^{1}(\Omega)}^{2} - \frac{1}{2\Lambda} \|b\|_{L^{\hat{p}}(\Omega)^{2}}^{2} \|z\|_{L^{\frac{2\hat{p}}{p-2}}(\Omega)}^{2}. \end{split}$$

Observe that the assumption $\hat{p} > 2$ implies that $2 < \frac{2\hat{p}}{\hat{p}-2} < \infty$. Now, we apply Lions' Lemma, [16, Chapter 2, Lemma 6.1], to the chain of embeddings $H^1(\Omega) \subset L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega) \subset L^2(\Omega)$, the first one being compact and the second one continuous, to deduce the existence of a constant C_0 depending on Λ , C_E and $\|b\|_{L^p(\Omega)^2}$ such that

$$\|z\|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega)} \leq \frac{\Lambda}{2^{3/2} \|b\|_{L^{\hat{p}}(\Omega)^{2}} C_{E}} \|z\|_{H^{1}(\Omega)} + C_{0} \|z\|_{L^{2}(\Omega)}.$$

From the last two inequalities we conclude (2.8) with

$$C_{\Lambda,E,b} = \frac{C_0^2 \|b\|_{L^{\hat{p}}(\Omega)^2}^2}{\Lambda}$$

and the proof is complete.

Remark 2.4. Notice that, to prove Lemma 2.3, we use neither $\nabla \cdot b \in L^{\hat{q}}(\Omega)$ nor $b \cdot n \in L^{\hat{q}}(\Gamma)$ for some $\hat{q} > 1$.

The adjoint operator of \mathcal{A} is \mathcal{A}^* . We have $\mathcal{A}^*z \in H^1(\Omega)'$ for every $z \in H^1(\Omega)$. In the next lemma, we justify that under the mild Assumption 2.1, we can integrate by parts and use the expected form of the adjoint Eq. (1.2b).

$$\langle \mathcal{A}^* z, y \rangle_{\Omega} = \int_{\Omega} \sum_{i,k=1}^2 a_{ki} \partial_{x_i} z \partial_{x_k} y \, \mathrm{d}x - \int_{\Omega} y \nabla \cdot (bz) \, \mathrm{d}x + \int_{\Gamma} y z b \cdot n \, \mathrm{d}x + \int_{\Omega} a_0 y z \, \mathrm{d}x.$$
 for all $y \in H^1(\Omega)$.

Proof. By definition

$$\langle \mathcal{A}^* z, y \rangle_{\Omega} = \langle \mathcal{A} y, z \rangle_{\Omega} \ \forall y, z \in H^1(\Omega)$$

and we only have to justify that, under Assumption 2.1, we can do integration by parts to get

$$\int_{\Omega} (b \cdot \nabla y) z \, \mathrm{d}x = -\int_{\Omega} y \nabla \cdot (bz) \, \mathrm{d}x + \int_{\Gamma} y z b \cdot n \, \mathrm{d}x.$$

This is equivalent to proving that we can apply the Gauss theorem to obtain

$$\int_{\Omega} \nabla \cdot (yzb) \, \mathrm{d}x = \int_{\Gamma} yzb \cdot n \, \mathrm{d}x.$$

Using that $y, z \in H^1(\Omega) \hookrightarrow L^r(\Omega)$ for all $r < +\infty$, $b \in L^{\hat{p}}(\Omega)^2$ for some $\hat{p} > 2$ and $\nabla \cdot b \in L^{\hat{q}}(\Omega)$ for some $\hat{q} > 1$, applying Hölder's inequality, we have

$$\nabla(yz) \cdot b = z\nabla y \cdot b + y\nabla z \cdot b \in L^{\frac{2\hat{p}}{2+\hat{p}}}(\Omega) \text{ and } yz\nabla \cdot b \in L^{\frac{\hat{q}+1}{2}}(\Omega),$$

so $\nabla \cdot (yzb) \in L^s(\Omega)$, where $s = \min\left\{\frac{2\hat{p}}{2+\hat{p}}, \frac{\hat{q}+1}{2}\right\}$ satisfies 1 < s < 2. From Assumption 2.1, it is also clear that $yzb \in L^s(\Omega)^2$, and using [17, Lema II.1.2.2], we deduce that yzb has a normal trace $yzb \cdot n$ defined in the space of $(W^{1-1/s',s'}(\Gamma))'$ via Gauss theorem: for every $v \in W^{1,s'}(\Omega)$

$$\langle yzb \cdot n, v \rangle_{(W^{1-1/s', s'}(\Gamma))', W^{1-1/s', s'}(\Gamma)} = \int_{\Omega} \nabla \cdot (vyzb) \, \mathrm{d}x.$$

Since we are assuming that $b \cdot n \in L^{\hat{q}}(\Gamma)$ for some $\hat{q} > 1$, then $yzb \cdot n \in L^{\frac{\hat{q}+1}{2}}(\Gamma) \hookrightarrow L^{s}(\Gamma)$. Therefore, we have that

$$\langle yzb \cdot n, v \rangle_{(W^{1-1/s',s'}(\Gamma))', W^{1-1/s',s'}(\Gamma)} = \int_{\Gamma} vyzb \cdot n \, \mathrm{d}x.$$

Taking v = 1 in the above equalities, we have that

$$\int_{\Omega} \nabla \cdot (yzb) \, \mathrm{d}x = \int_{\Gamma} yzb \cdot n \, \mathrm{d}x$$

and the proof is complete.

Next, we adapt the proof of [8, Theorem 2.2] to show existence and uniqueness of the solution of the state equation.

Lemma 2.6. Under Assumption 2.1, the linear operator $\mathcal{A} : H^1(\Omega) \longrightarrow H^1(\Omega)'$ is an isomorphism.

 \Box

Proof. Let us first see that \mathcal{A} is injective. Consider $y \in H^1(\Omega)$ such that $\mathcal{A}y = 0$. We will prove that $y \leq 0$. The contrary inequality follows by arguing on -y. Suppose there exists some $\mathcal{O} \subset \Omega$ with positive measure such that y(x) > 0 if $x \in \mathcal{O}$. Take $0 < \rho < \operatorname{ess\,sup}_{x \in \Omega} y(x) \leq +\infty$ and define $y_{\rho}(x) = (y(x) - \rho)^+ = \max\{y(x) - \rho, 0\}$. Denote $\Omega_{\rho} = \{x \in \Omega : \nabla y_{\rho}(x) \neq 0\}$. Notice that $y_{\rho} \in H^1(\Omega)$,

$$\nabla y_{\rho}(x) = \begin{cases} \nabla y(x) & \text{if } y(x) > \rho \\ 0 & \text{if } y(x) \le \rho, \end{cases}$$

which means that $\Omega_{\rho} \subset \{x : y(x) > \rho\}$. We also remark that $y_{\rho}(x) = 0$ if $y(x) \leq 0$, and that $y(x) \geq y_{\rho}(x) \geq 0$ if $y(x) \geq 0$. Using these properties, and Hölder's and Young's inequalities, we have that

$$\begin{aligned} 0 &= \langle \mathcal{A}y, y_{\rho} \rangle_{\Omega} = \int_{\Omega} \sum_{i,k=1}^{2} a_{ik} \partial_{x_{i}} y \partial_{x_{k}} y_{\rho} \, \mathrm{d}x + \int_{\Omega} (b \cdot \nabla y) y_{\rho} \, \mathrm{d}x + \int_{\Omega} a_{0} y y_{\rho} \, \mathrm{d}x \\ &\geq \int_{\Omega_{\rho}} \sum_{i,k=1}^{2} a_{ik} \partial_{x_{i}} y_{\rho} \partial_{x_{k}} y_{\rho} \, \mathrm{d}x + \int_{\Omega_{\rho}} (b \cdot \nabla y_{\rho}) y_{\rho} \, \mathrm{d}x + \int_{\Omega} a_{0} y_{\rho} y_{\rho} \, \mathrm{d}x \\ &\geq \Lambda \| \nabla y_{\rho} \|_{L^{2}(\Omega_{\rho})}^{2} - \| b \|_{L^{\hat{p}}(\Omega_{\rho})^{2}} \| \nabla y_{\rho} \|_{L^{2}(\Omega_{\rho})} \| y_{\rho} \|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega_{\rho})}^{2} + \frac{\Lambda}{2} \| y_{\rho} \|_{L^{2}(E)}^{2} \\ &\geq \frac{\Lambda}{2} \| \nabla y_{\rho} \|_{L^{2}(\Omega_{\rho})}^{2} - \frac{1}{2\Lambda} \| b \|_{L^{\hat{p}}(\Omega_{\rho})^{2}}^{2} \| y_{\rho} \|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega_{\rho})}^{2} + \frac{\Lambda}{2} \| y_{\rho} \|_{L^{2}(E)}^{2} \\ &= \frac{\Lambda}{2} \| \nabla y_{\rho} \|_{L^{2}(\Omega)}^{2} - \frac{1}{2\Lambda} \| b \|_{L^{\hat{p}}(\Omega_{\rho})^{2}}^{2} \| y_{\rho} \|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega_{\rho})}^{2} + \frac{\Lambda}{2} \| y_{\rho} \|_{L^{2}(E)}^{2} \end{aligned}$$

Next we use that $\Omega_{\rho} \subset \Omega$, (2.3), (2.2) and the just proved inequality to obtain:

$$\begin{split} \|y_{\rho}\|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega_{\rho})}^{2} &\leq \|y_{\rho}\|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega)}^{2} \leq K_{\Omega,\frac{2p}{p-2}}^{2}\|y_{\rho}\|_{H^{1}(\Omega)}^{2} \\ &\leq 2K_{\Omega,\frac{2\hat{p}}{\hat{p}-2}}^{2}C_{E}^{2}\left(\|\nabla y_{\rho}\|_{L^{2}(\Omega)}^{2}+\|y_{\rho}\|_{L^{2}(E)}^{2}\right) \\ &\leq \frac{2K_{\Omega,\frac{2\hat{p}}{\hat{p}-2}}^{2}C_{E}^{2}}{\Lambda^{2}}\|b\|_{L^{\hat{p}}(\Omega_{\rho})^{2}}^{2}\|y_{\rho}\|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega_{\rho})}^{2} \end{split}$$

We can deduce from this a positive lower bound for the norm of b in $L^{\hat{p}}(\Omega_{\rho})^2$ independent of ρ :

$$\|b\|_{L^{\hat{p}}(\Omega_{\rho})^{2}} \ge \frac{\Lambda}{\sqrt{2}K_{\Omega,\frac{2\hat{p}}{\hat{p}-2}}C_{E}} > 0.$$

But we have that $|\Omega_{\rho}| \to 0$ as $\rho \to \operatorname{ess\,sup}_{x \in \Omega} y(x)$; see [8, Theorem 2.2]. So we have achieved a contradiction.

Finally we have just to check that the range of \mathcal{A} is dense and closed. Since we already have established Gårding's inequality (2.8) for the operator \mathcal{A} , the proof of closeness done in [8, Theorem 2.2] applies to our case changing the norms in $H_0^1(\Omega)$ and $H^{-1}(\Omega)$ respectively by the norms in $H^1(\Omega)$ and its dual space, and thus it is omitted. By a well known duality argument, the denseness of the range of \mathcal{A} follows from the injectivity of \mathcal{A}^* .

The argument used above to obtain the injectivity of \mathcal{A} does not work for \mathcal{A}^* . Notice that at one moment we use that $\int_{\Omega} (b \cdot \nabla y) y_{\rho} dx = \int_{\Omega_{\rho}} (b \cdot \nabla y_{\rho}) y_{\rho} dx$. When dealing with the adjoint operator, we would find the term $\int_{\Omega} (b \cdot \nabla z_{\rho}) z \, dx$, which in general is different from $\int_{\Omega_{\rho}} (b \cdot \nabla z_{\rho}) z_{\rho} \, dx$. But we can obtain injectivity of the adjoint operator as follows. Consider $z \in H^1(\Omega)$ such that $\mathcal{A}^* z = 0$. For all $\varepsilon \geq 0$ define

$$\Omega^{\varepsilon} = \{ x \in \Omega : |z(x)| > \varepsilon \}.$$

Let us see that $|\Omega^0| = 0$, which readily implies that z = 0. Let us define $z^{\varepsilon}(x) = \operatorname{Proj}_{[-\varepsilon,\varepsilon]}(z(x))$. Using integration by parts, that z = 0 in $\Omega \setminus \Omega^0$, that $\nabla z^{\varepsilon} = 0$ in Ω^{ε} and $\nabla z^{\varepsilon} = \nabla z$ in $\Omega \setminus \Omega^{\varepsilon}$, and that $zz^{\varepsilon} \ge (z^{\varepsilon})^2$, we have

$$\begin{split} 0 &= \langle \mathcal{A}^* z, z^{\varepsilon} \rangle_{\Omega} \\ &= \int_{\Omega} \sum_{i,k=1}^{2} a_{ki} \partial_{x_{i}} z \partial_{x_{k}} z^{\varepsilon} \, \mathrm{d}x - \int_{\Omega} z^{\varepsilon} \nabla \cdot (bz) \, \mathrm{d}x + \int_{\Gamma} z^{\varepsilon} z b \cdot n \, \mathrm{d}x + \int_{\Omega} a_{0} z^{\varepsilon} z \, \mathrm{d}x \\ &= \int_{\Omega} \sum_{i,k=1}^{2} a_{ki} \partial_{x_{i}} z \partial_{x_{k}} z^{\varepsilon} \, \mathrm{d}x + \int_{\Omega} z b \cdot \nabla z^{\varepsilon} \, \mathrm{d}x + \int_{\Omega} a_{0} z^{\varepsilon} z \, \mathrm{d}x \\ &\geq \Lambda \| \nabla z^{\varepsilon} \|_{L^{2}(\Omega)^{2}} - \| b \|_{L^{\hat{p}}(\Omega^{0} \setminus \Omega^{\varepsilon})} \| \nabla z^{\varepsilon} \|_{L^{2}(\Omega)^{2}} \| z^{\varepsilon} \|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega^{0} \setminus \Omega^{\varepsilon})} + \frac{\Lambda}{2} \| z^{\varepsilon} \|_{L^{2}(E)}^{2} \\ &\geq \frac{\Lambda}{2} \| \nabla z^{\varepsilon} \|_{L^{2}(\Omega)^{2}} - \frac{1}{2\Lambda} \| b \|_{L^{\hat{p}}(\Omega^{0} \setminus \Omega^{\varepsilon})}^{2} \| z^{\varepsilon} \|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega^{0} \setminus \Omega^{\varepsilon})} + \frac{\Lambda}{2} \| z^{\varepsilon} \|_{L^{2}(E)}^{2}. \end{split}$$

So, using this and that $|z^{\varepsilon}(x)| \leq \varepsilon$ for a.e. $x \notin \Omega^{\varepsilon}$ we get $||z^{\varepsilon}||^{2}_{H^{1}(\Omega)} \leq 2C_{E}^{2} \left(||\nabla z^{\varepsilon}||_{L^{2}(\Omega)^{2}} + ||z^{\varepsilon}||^{2}_{L^{2}(E)} \right)$

$$\leq \frac{2C_E^2}{\Lambda^2} \|b\|_{L^{\hat{p}}(\Omega^0 \setminus \Omega^{\varepsilon})}^2 \|z^{\varepsilon}\|_{L^{\frac{2\hat{p}}{\hat{p}-2}}(\Omega^0 \setminus \Omega^{\varepsilon})}^2 \leq \frac{2C_E^2}{\Lambda^2} \|b\|_{L^{\hat{p}}(\Omega^0 \setminus \Omega^{\varepsilon})}^2 |\Omega^0 \setminus \Omega^{\varepsilon}|^{\frac{\hat{p}-2}{\hat{p}}} \varepsilon^2.$$

On the other hand, using that $|z^{\varepsilon}| = \varepsilon$ in Ω^{ε} and the previous inequality, we have

$$\begin{aligned} |\Omega^{\varepsilon}| &= \frac{1}{\varepsilon^2} \int_{\Omega^{\varepsilon}} z^{\varepsilon}(x)^2 \, \mathrm{d}x \le \frac{1}{\varepsilon^2} \|z^{\varepsilon}\|_{L^2(\Omega)}^2 \le \frac{1}{\varepsilon^2} \|z^{\varepsilon}\|_{H^1(\Omega)}^2 \\ &\le \frac{2C_E^2}{\Lambda^2} \|b\|_{L^{\hat{p}}(\Omega^0 \setminus \Omega^{\varepsilon})}^2 |\Omega^0 \setminus \Omega^{\varepsilon}|^{\frac{\hat{p}-2}{\hat{p}}}. \end{aligned}$$

Since $|\Omega^0 \setminus \Omega^{\varepsilon}| = \max\{x \in \Omega : 0 < |z(x)| < \varepsilon\} \to 0$ as $\varepsilon \to 0$, we have proved that $|\Omega^{\varepsilon}| \to 0$ as $\varepsilon \to 0$ and hence $|\Omega^0| = 0$.

Corollary 2.7. Under Assumption 2.1, the linear operator \mathcal{A}^* : $H^1(\Omega) \longrightarrow H^1(\Omega)'$ is an isomorphism.

3. Regularity of the Solution of the State and Adjoint State Equations

To obtain further regularity, from now on we will suppose

Assumption 3.1. The coefficients a_{ik} belong to $C^{0,1}(\overline{\Omega}), 1 \leq i, k, \leq 2$.

Let us denote by m the number of sides of Γ and $\{S_j\}_{j=1}^m$ its vertices, ordered counterclockwise. For convenience denote also $S_0 = S_m$ and $S_{m+1} = S_1$. We denote by Γ_j the side of Γ connecting S_j and S_{j+1} , and by $\omega_j \in (0, 2\pi)$ the angle interior to Ω at S_j , i.e., the angle defined by Γ_j and Γ_{j-1} , measured counterclockwise. Notice that $\Gamma_0 = \Gamma_m$. We use (r_j, θ_j) as local polar coordinates at S_j , with $r_j = |x - S_j|$ and θ_j the angle defined by Γ_j and the segment $[S_j, x]$. In order to describe and analyze the regularity of the functions near the corners, we will introduce for every $j \in \{1, \ldots, m\}$ the infinite cone

$$K_j = \{ x \in \mathbb{R}^2 : 0 < r_j, \ 0 < \theta_j < \omega_j \}.$$

For every $j \in \{1, ..., m\}$ we call A_j the operator with constant coefficients, corresponding to the corner S_j , given by

$$A_j y = \sum_{i,k=1}^2 \partial_{x_k} (a_{ik}(S_j) \partial_{x_i} y).$$

We denote by λ_j the leading singular exponent associated with the operator A_j at the corner S_j , i.e., the smallest $\lambda_j > 0$ such that there exists a solution of the form $y_j = r_j^{\lambda_j} \varphi_j(\theta_j)$, with φ_j smooth enough, of

$$A_j y_j = 0$$
 in K_j , $\partial_{n_{A_j}} y_j = 0$ on ∂K_j .

For instance, for $Ay = -\Delta y$ we have $\lambda_j = \pi/\omega_j$. We denote $\lambda = \min\{\lambda_j\}$.

With the usual technique of taking a partition of the unity to localize the problem in the corners, freezing the coefficients and doing an appropriate linear change of variable, the classical results for the Laplace operator are also valid in our case; see, e.g. [18, Section 2.1] for a detailed example of application of this technique. Notice that the symmetry hypothesis $a_{ik} = a_{ki}$ introduced in Assumption 2.1 implies that the same change of variable that transforms A_j into $-\Delta$ will transform the conormal derivative $\partial_{n_{A_j}}$ into the normal derivative ∂_n in the new variables, and not in an oblique derivative.

We continue with regularity results for problems with $b \equiv 0$ and $a_0 \equiv 0$ and use the standard Sobolev and Sobolev–Slobodetskiĭ spaces but also weighted Sobolev spaces as follows. Let $k \in \mathbb{N}_0$, $1 \leq p \leq \infty$, and $\vec{\beta} = (\beta_1, \ldots, \beta_m)^T \in \mathbb{R}^m$, $j \in \{1, \ldots, m\}$. For ball-neighborhoods Ω_{R_j} of S_j with

radius $R_j \leq 1$ and $\Omega^0 := \Omega \setminus \bigcup_{j=1}^m \Omega_{R_j/2}$ we define norms via

$$\begin{split} \|v\|_{W^{k,p}_{\beta_{j}}(\Omega_{R_{j}})}^{p} &= \sum_{|\alpha| \leq k} \|r_{j}^{\beta_{j}} D^{\alpha} v\|_{L^{p}(\Omega_{R_{j}})}^{p}, \\ \|v\|_{V^{k,p}_{\beta_{j}}(\Omega_{R_{j}})}^{p} &= \sum_{|\alpha| \leq k} \|r_{j}^{\beta_{j}-k+|\alpha|} D^{\alpha} v\|_{L^{p}(\Omega_{R_{j}})}^{p}, \end{split}$$

where the standard modification for $p = \infty$ is used. The spaces $W^{k,p}_{\vec{\beta}}(\Omega)$ and $V^{k,p}_{\vec{\beta}}(\Omega)$ denote the set of all functions v such that

$$\begin{aligned} \|v\|_{W^{k,p}_{\vec{\beta}}(\Omega)} &:= \|v\|_{W^{k,p}(\Omega^0)} + \sum_{j=1}^m \|v\|_{W^{k,p}_{\vec{\beta}_j}(\Omega_{R_j})}, \\ \|v\|_{V^{k,p}_{\vec{\beta}}(\Omega)} &:= \|v\|_{W^{k,p}(\Omega^0)} + \sum_{j=1}^m \|v\|_{V^{k,p}_{\vec{\beta}_j}(\Omega_{R_j})}, \end{aligned}$$

respectively, are finite. The corresponding seminorms are defined by setting $|\alpha| = k$ instead of $|\alpha| \leq k$. For the definition of the corresponding trace spaces $W_{\vec{\beta}}^{k-1/p,p}(\Gamma_j), \ V_{\vec{\beta}}^{k-1/p,p}(\Gamma_j), \ W_{\vec{\beta}}^{k-1/p,p}(\Gamma)$ and $V_{\vec{\beta}}^{k-1/p,p}(\Gamma)$ we refer to [13, Sect. 6.2.10], see also [19, Section 2.2]. We will also use the notation $L_{\vec{\beta}}^p(\Omega)$ for $W_{\vec{\beta}}^{0,p}(\Omega)$.

Lemma 3.2. Suppose that Assumption 3.1 holds. Consider $f \in H^1(\Omega)'$ and $g \in H^{1/2}(\Gamma)'$ such that

$$\langle f, 1 \rangle_{\Omega} + \langle g, 1 \rangle_{\Gamma} = 0,$$

and let $y \in H^1(\Omega)$ be the unique solution, up to a constant, of

$$\int_{\Omega} \sum_{i,k=1}^{2} a_{ik} \partial_{x_i} y \partial_{x_k} z \, \mathrm{d}x = \langle f, z \rangle_{\Omega} + \langle g, z \rangle_{\Gamma} \quad \forall z \in H^1(\Omega)$$

We have the following regularity results.

(a) If
$$f \in H^{2-t}(\Omega)'$$
, and $g \in \prod_{j=1}^{m} H^{t-3/2}(\Gamma_j)$ for some $1 < t < 1 + \lambda, t \le 2$,
then

then

$$y \in H^t(\Omega) \text{ and } |y|_{H^t(\Omega)} \le C_{A,t} \Big(||f||_{H^{2-t}(\Omega)'} + \sum_{j=1}^m ||g||_{H^{t-3/2}(\Gamma_j)} \Big).$$

(b) If
$$f \in L^{r}(\Omega)$$
 and $g \in \prod_{j=1}^{m} W^{1-1/r,r}(\Gamma_{j})$ for some $1 < r < \frac{2}{2-\lambda}$ if $\lambda < 2$,
 $r > 1$ arbitrary if $\lambda \ge 2$, then

$$y \in W^{2,r}(\Omega) \text{ and } |y|_{W^{2,r}(\Omega)} \le C_{A,r} \Big(||f||_{L^{r}(\Omega)} + \sum_{j=1}^{m} ||g||_{W^{1-1/r,r}(\Gamma_{j})} \Big).$$

(c) Consider $s \in (1,\infty)$ and $\vec{\beta}$ such that $2 - \frac{2}{s} - \lambda_j < \beta_j < 2 - \frac{2}{s}, \ \beta_j \ge 0$ for all $j \in \{1,\ldots,m\}$. If $f \in L^s_{\vec{\beta}}(\Omega)$ and $g \in \prod_{j=1}^m V^{1-1/s,s}_{\vec{\beta}}(\Gamma_j)$, then

$$y \in W^{2,s}_{\vec{\beta}}(\Omega) \text{ and } |y|_{W^{2,s}_{\vec{\beta}}(\Omega)} \leq C_{A,\vec{\beta}} \Big(\|f\|_{L^{s}_{\vec{\beta}}(\Omega)} + \sum_{j=1}^{m} \|g\|_{V^{1-1/s,s}_{\vec{\beta}}(\Gamma_{j})} \Big).$$

Remark 3.3. Let us briefly comment on the function spaces appearing in the lemma. Notice that for t = 2, $H^{t-2}(\Omega) = H^{2-t}(\Omega) = L^2(\Omega)$, and for 3/2 < t, $H^{2-t}(\Omega) = H_0^{2-t}(\Omega)$ and hence $H^{t-2}(\Omega) = H^{2-t}(\Omega)'$. Nevertheless, for 1 < t < 3/2, $H^{2-t}(\Omega)' \neq H^{t-2}(\Omega)$. Also take into account that

$$\prod_{j=1}^{m} H^{t-3/2}(\Gamma_j) = H^{t-3/2}(\Gamma) \text{ if } t < 2, \quad \prod_{j=1}^{m} W^{1-1/r,r}(\Gamma_j) = W^{1-1/r,r}(\Gamma) \text{ if } r < 2.$$

We remark that the mapping $y \mapsto \partial_{n_A} y$ is linear and continuous from $H^2(\Omega)$ onto $\prod_{j=1}^m H^{1/2}(\Gamma_j)$; see [10, Theorem 1.5.2.8].

Regarding weighted spaces, we notice that $V_{\vec{\beta}}^{1-1/s,s}(\Gamma) = W_{\vec{\beta}}^{1-1/s,s}(\Gamma)$ if $\beta_j > 1 - \frac{2}{s}$ or $\beta_j < -\frac{2}{s}$ for all $j \in \{1, \ldots, m\}$, while these spaces differ by a constant in the vicinity of each corner S_j where $-\frac{2}{s} < \beta_j < 1 - \frac{2}{s}$, see [12, Theorem 2.1] or [14, page 131].

Proof of Lemma 3.2. The result in (a) can be deduced from [20, Theorem 9.2] for 1 < t < 3/2, from [21, Theorem (23.3)] for 3/2 < t < 2, and from [10, Corollary 4.4.4.14] for t = 2. The case t = 3/2 follows by interpolation. Statement (b) follows from [10, Corollary 4.4.4.14]. Part (c) follows by standard arguments but we did not find this particular result in the literature. Therefore we sketch the proof here for the case of constant coefficients. As said above, the result in the case of Lipschitz coefficients follows from this one using the localization-and-freezing technique.

We will use [13, Theorem 1.2.5] stating a similar result for a cone K and weighted V-spaces. For the problem under consideration and in our notation it says that $y \in V_{\beta}^{2,s}(K)$ if $f \in L_{\beta}^{s}(K)$ and $g \in V_{\beta}^{1-1/s,s}(\partial K \setminus O)$ provided that $s \in (1, \infty)$ and $2 - \frac{2}{s} - \beta \notin \{k\lambda, k \in \mathbb{Z}\}$. To satisfy the latter condition we assume $2 - \frac{2}{s} - \lambda_{j} < \beta_{j} < 2 - \frac{2}{s}$ for all $j \in \{1, \ldots, m\}$.

The reformulation from the vicinity of a vertex of the domain Ω is achieved by using cut-off functions $\zeta_j : \Omega \to [0,1]$ with $\zeta_j \equiv 1$ in $\Omega_{R_j/2}$, $\zeta_j \equiv 0$ in $\Omega \setminus \Omega_{R_j}$, and $\partial_{n_{A_j}} \zeta_j = 0$ on $\partial \Omega \cap \partial \Omega_{R_j}$. We split $y \in H^1(\Omega)$ into

$$y = \sum_{j=1}^{m} y_j + w$$
, where $y_j = \zeta_j (y - y(S_j))$.

With this construction we get $y_j(S_j) = 0$ and $\operatorname{supp} y_j = \overline{\Omega}_{R_j}$ such that we can consider the problem $Ay_j = f_j$ with Neumann boundary condition $\partial_{n_{A_j}} y_j =$

 $g_j = \zeta_j g$ in the cone K_j . For f_j , we have

$$f_j = A(\zeta_j(y - y(S_j))) = \begin{cases} \zeta_j f & \text{in } \Omega_{R_j/2} \\ \zeta_j f - \mathfrak{b}_j \cdot \nabla y - \mathfrak{a}_j(y - y(S_j)) & \text{in } \Omega_{R_j} \setminus \Omega_{R_j/2} \\ 0 & \text{in } K_j \setminus \Omega_{R_j} \end{cases}$$

with smooth functions \mathfrak{b}_j and \mathfrak{a}_j due to the constant coefficients in A. From $f \in L^s_{\overline{\beta}}(\Omega)$ and $y \in H^1(\Omega)$ we conclude $f_j \in L^{\hat{s}}_{\beta_j}(K)$, $\hat{s} = \min(s, 2)$ where we use that $\beta_j \geq 0$. Moreover, the assumption $g \in \prod_{j=1}^m V^{1-1/s,s}_{\overline{\beta}}(\Gamma_j)$ leads to $g_j \in V^{1-1/s,s}_{\beta_j}(\partial K_j \setminus O_j)$ such that [13, Theorem 1.2.5] leads to $y_j \in V^{2,\hat{s}}_{\beta_j}(K_j) \hookrightarrow W^{2,\hat{s}}_{\beta_j}(K_j)$. Since the function w does not contain corner singularities, hence $w \in W^{2,s}(\Omega)$, we obtain $y \in W^{2,\hat{s}}_{\overline{\beta}}(\Omega)$. If $s \leq 2$ we are done.

Otherwise, when s > 2, we have $y \in H^2(\Omega_{R_j} \setminus \Omega_{R_j/2}) \hookrightarrow W^{1,s}(\Omega_{R_j} \setminus \Omega_{R_j/2})$, and we reiterate $f_j \in L^s_{\beta_j}(K)$ and $y_j \in V^{2,s}_{\beta_j}(K_j) \hookrightarrow W^{2,\hat{s}}_{\beta_j}(K_j)$ leading to $y \in W^{2,s}_{\vec{\beta}}(\Omega)$.

Theorem 3.4. Suppose that Assumptions 2.1 and 3.1 hold. Consider $f \in H^1(\Omega)'$ and $u \in H^{1/2}(\Gamma)'$ and let $y \in H^1(\Omega)$ be the unique solution of

$$\langle \mathcal{A}y, z \rangle_{\Omega} = \langle f, z \rangle_{\Omega} + \langle u, z \rangle_{\Gamma} \quad \forall z \in H^{1}(\Omega).$$
 (3.1)

We have the following regularity results.

(a) If $a_0 \in L^q(\Omega)$, $f \in H^{2-t}(\Omega)'$ and $u \in \prod_{j=1}^m H^{t-3/2}(\Gamma_j)$ for some t such that $1 < t < 1 + \lambda$, $t \leq 2$ and $q = \frac{2}{3-t}$, then $y \in H^t(\Omega)$ and there exists a constant $C_{\mathcal{A},t} > 0$ such that

$$\|y\|_{H^t(\Omega)} \le C_{\mathcal{A},t}(\|f\|_{H^{2-t}(\Omega)'} + \sum_{j=1}^m \|u\|_{H^{t-3/2}(\Gamma_j)}).$$

(b) If $a_0 \in L^r(\Omega)$, $f \in L^r(\Omega)$ and $u \in \prod_{j=1}^m W^{1-1/r,r}(\Gamma_j)$ for some $r \in (1,\hat{p}]$ satisfying $r < \frac{2}{2-\lambda}$ in case of $\lambda < 2$, then $y \in W^{2,r}(\Omega)$ and there exists a constant $C_{\mathcal{A},r} > 0$ such that

$$\|y\|_{W^{2,r}(\Omega)} \le C_{\mathcal{A},r}(\|f\|_{L^{r}(\Omega)} + \sum_{j=1}^{m} \|u\|_{W^{1-1/r,r}(\Gamma_{j})})$$

(c) If $a_0 \in L^p_{\vec{\beta}}(\Omega)$, $f \in L^p_{\vec{\beta}}(\Omega)$ and $u \in \prod_{j=1}^m W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)$ for some $p \in (1,2]$ and some $\vec{\beta}$ such that $2 - \frac{2}{p} - \lambda_j < \beta_j < 2 - \frac{2}{p}$ and $\beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$, then $y \in W^{2,p}_{\vec{\beta}}(\Omega)$ and there exists a constant $C_{\mathcal{A},\vec{\beta},p} > 0$ such that

$$\|y\|_{W^{2,p}_{\vec{\beta}}(\Omega)} \le C_{\mathcal{A},\vec{\beta},p}(\|f\|_{L^{p}_{\vec{\beta}}(\Omega)} + \sum_{j=1}^{m} \|u\|_{W^{1-1/p,p}_{\vec{\beta}}(\Gamma_{j})}).$$

Proof. Let us define

$$F = -b \cdot \nabla y - a_0 y.$$

From the proof of Lemma 2.3, we know that $F \in H^1(\Omega)'$. Also, taking z = 1 in (3.1), we have that

$$\langle f + F, 1 \rangle_{\Omega} + \langle u, 1 \rangle_{\Gamma} = 0$$

so the conditions of Lemma 3.2 apply to the problem

$$\langle Ay, z \rangle_{\Omega} = \langle f + F, z \rangle_{\Omega} + \langle u, z \rangle_{\Gamma} \quad \forall z \in H^{1}(\Omega).$$

We have to investigate the regularity of F.

(a) For $1 < \tau \leq t$, define $S = \{z \in H^{2-\tau}(\Omega) : \|z\|_{H^{2-\tau}(\Omega)} = 1\}$. We have that $F \in H^{2-\tau}(\Omega)'$ if and only if

$$\|F\|_{H^{2-\tau}(\Omega)'} = \sup_{z \in S} |\langle F, z \rangle_{\Omega}| < +\infty.$$

Applying Hölder's inequality, we can deduce the existence of a constant $C_{\Omega} > 0$, that may depend on the measure of Ω , such that

$$\begin{aligned} |\langle F, z \rangle_{\Omega}| &= \left| \int_{\Omega} (b \cdot \nabla y + a_0 y) z \, \mathrm{d}x \right| \\ &\leq C_{\Omega} \big(\|b\|_{L^{\hat{p}}(\Omega)^2} \|\nabla y\|_{L^{r_p}(\Omega)} + \|a_0\|_{L^q(\Omega)} \|y\|_{L^{r_q}(\Omega)} \big) \|z\|_{L^s(\Omega)} \quad (3.2) \end{aligned}$$

where

$$\frac{1}{\hat{p}} + \frac{1}{r_p} + \frac{1}{s} \le 1, \quad \frac{1}{q} + \frac{1}{r_q} + \frac{1}{s} \le 1.$$
(3.3)

Let us also notice that $H^{2-\tau}(\Omega) \hookrightarrow L^s(\Omega)$ if and only if

$$\tau = 1 + \frac{2}{s}.\tag{3.4}$$

We will apply a boot-strap argument.

Step 1. We know that $y \in H^1(\Omega)$, so $r_p = 2$ and for r_q we can take any real number. Noting that q > 1, using (3.3) and taking

$$\frac{1}{s} = \min\left\{1 - \frac{1}{\hat{p}} - \frac{1}{r_p}, 1 - \frac{1}{q} - \frac{1}{r_q}\right\},\,$$

we have that 1/s > 0 for r_q big enough and both conditions in (3.3) are satisfied. Hence we deduce that $F \in H^{2-\tau}(\Omega)'$. Since $u \in \prod_{j=1}^m H^{t-3/2}(\Gamma_j)$, a direct application of Lemma 3.2 yields that $y \in H^{\min\{t,\tau\}}(\Omega)$. If $\tau \ge t$, the proof is complete. Step 2. Otherwise we have that $\nabla y \in H^{\tau-1}(\Omega)^2 \hookrightarrow L^{r_p}(\Omega)^2$ for $1 \qquad \tau$

$$\frac{1}{r_p} = 1 - \frac{7}{2}$$

and, since $\tau > 1$, we can take $r_q = +\infty$. As before, we select

$$\frac{1}{s} = \min\left\{1 - \frac{1}{\hat{p}} - \frac{1}{r_p}, 1 - \frac{1}{q}\right\}.$$

We have two possibilities now.

Step 3. If $\frac{1}{s} = 1 - \frac{1}{q}$, then, applying (3.4) and taking into account our choice of q, we have that $y \in H^{\hat{\tau}}(\Omega)$ with

$$\hat{\tau} = 1 + \frac{2}{s} = 3 - \frac{2}{q} = t,$$

and the proof is complete.

Step 4. Otherwise,
$$\frac{1}{s} = 1 - \frac{1}{\hat{p}} - \frac{1}{r_p}$$
 and we will have $y \in H^{\hat{\tau}}(\Omega)$ with

$$\hat{\tau} = 1 + \frac{2}{s} = 1 + 2 - \frac{2}{\hat{p}} - (2 - \tau) = \tau + 1 - \frac{2}{\hat{p}},$$

and we have advanced a fixed amount $1 - \frac{2}{\hat{p}}$. If $\hat{\tau} \ge t$, the proof is complete. Step 5. In other case, we can redefine $\tau = \hat{\tau}$ and go back to step 2.

Every time we repeat the process, either we finish the proof or we increment the size of τ by the fixed amount $1 - \frac{2}{\hat{p}}$, so the proof will end in a finite number of steps.

(b) From the Sobolev embedding theorem, we have that

$$f \in L^r(\Omega) \hookrightarrow H^{2-t}(\Omega)'$$
 and $u \in \prod_{j=1}^m W^{1-1/r,r}(\Gamma_j) \hookrightarrow \prod_{j=1}^m H^{t-3/2}(\Gamma_j)$

for $t = \min\{2, 3-2/r\}$. The conditions imposed on r imply that $1 < t < 1+\lambda$, $t \leq 2$, so we can apply Theorem 3.4(a) to obtain $y \in H^t(\Omega)$ and we readily have that $y \in L^{\infty}(\Omega)$ and hence $a_0 y \in L^r(\Omega)$. Let us investigate the regularity of $b \cdot \nabla y$. We use again a boot-strap argument.

We have that $\nabla y \in H^{t-1}(\Omega) \hookrightarrow L^{\frac{2}{2-t}}(\Omega)$. Therefore $b \cdot \nabla y \in L^s(\Omega)$ where

$$\frac{1}{s} = \frac{1}{\hat{p}} + \frac{2-t}{2} < 1.$$

Applying Lemma 3.2(b), we have that $y \in W^{2,\min\{s,r\}}(\Omega)$. If $s \ge r$, the proof is complete. Otherwise, we have that $\nabla y \in W^{1,s}(\Omega) \hookrightarrow L^{s^*}(\Omega)$, with

$$\frac{1}{s^*} = \frac{1}{s} - \frac{1}{2}$$

Therefore $b \cdot \nabla y \in L^{\hat{s}}(\Omega)$ where

$$\frac{1}{\hat{s}} = \frac{1}{\hat{p}} + \frac{1}{s^*} = \frac{1}{\hat{p}} + \frac{1}{s} - \frac{1}{2} = \frac{1}{s} - \left(\frac{1}{2} - \frac{1}{\hat{p}}\right).$$

If $\frac{1}{\hat{s}} \leq \frac{1}{r}$, then the proof is complete. Otherwise, we can rename $s := \hat{s}$ and repeat the argument subtracting at each step the positive constant $\frac{1}{2} - \frac{1}{\hat{p}}$ until 1 = 1

 $\frac{1}{\hat{s}} \le \frac{1}{r}.$

(c) To obtain this result, we want to apply Lemma 3.2(c), but the boundary datum in that result is in the space $\prod_{j=1}^{m} V_{\vec{\beta}}^{1-1/p,p}(\Gamma_j)$, while the boundary datum in this result is in $\prod_{j=1}^{m} W_{\vec{\beta}}^{1-1/p,p}(\Gamma_j)$. Taking into account Remark 3.3, it is clear that for p < 2, the condition $\beta_j \ge 0$ implies that $\beta_j > 1 - 2/p$ and hence $W_{\vec{\beta}}^{1-1/p,p}(\Gamma_j) = V_{\vec{\beta}}^{1-1/p,p}(\Gamma_j)$ for all $j \in \{1, \ldots, m\}$. If p = 2, we define

$$u_s = \sum_{\beta_j > 0} u\zeta_j,$$

where the ζ_j are the cutoff functions introduced in the proof of Lemma 3.2(c). Taking into account again Remark 3.3 and noting that $u_s \equiv 0$ in a neighbourhood of the corners S_j with $\beta_j = 0$, it is readily deduced that $u_s \in \prod_{j=1}^m V_{\beta}^{1-1/p,p}(\Gamma_j)$. We also have that the function $u_r = u - u_s \in \prod_{j=1}^m W^{1-1/p,p}(\Gamma_j)$, because $u_r \equiv 0$ in a neighbourhood of the corners S_j such that $\beta_j > 0$. In the same way we define

$$f_s = \sum_{\beta_j > 0} f\zeta_j \in L^p_{\vec{\beta}}(\Omega) \text{ and } f_r = f - f_s \in L^p(\Omega),$$

and consider $y_s, y_r \in H^1(\Omega)$ such that

 $\langle \mathcal{A}y_r, z \rangle_{\Omega} = \langle f_r, z \rangle_{\Omega} + \langle u_r, z \rangle_{\Gamma}, \text{ and } \langle \mathcal{A}y_s, z \rangle_{\Omega} = \langle f_s, z \rangle_{\Omega} + \langle u_s, z \rangle_{\Gamma} \quad \forall z \in H^1(\Omega),$

so that $y = y_r + y_s$. As an application of Theorem 3.4(b), $y_r \in W^{2,2}(\Omega)$, which is continuously embedded in $W^{2,2}_{\vec{\beta}}(\Omega)$ because $\beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$.

Taking into account the above considerations, in the rest of the proof we assume that $\beta_j > 1 - 2/p$. If p < 2 then this holds, as discussed before. If p = 2, we denote $u_s = u$ to treat both cases simultaneously, and hence we can use both that $u \in \prod_{j=1}^m W_{\vec{\beta}}^{1-1/p,p}(\Gamma_j)$, which is needed to have an embedding in a non-weighted Sobolev space, and $u \in \prod_{j=1}^m V_{\vec{\beta}}^{1-1/p,p}(\Gamma_j)$, which is needed to apply Lemma 3.2(c).

From [19, Lemma 2.29(ii)], we deduce that $L^p_{\vec{\beta}}(\Omega) \hookrightarrow L^r(\Omega)$ for all $r < \frac{2}{\beta_j + 2/p} \leq \frac{2}{2/p} = p$ for all $j \in \{1, \ldots, m\}$. On the other hand, using the definition of the $\prod_{j=1}^m W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)$ -norm and [19, Lemma 2.29(i)], we have

the embedding $\prod_{j=1}^{m} W_{\vec{\beta}}^{1-1/p,p}(\Gamma_j) \hookrightarrow \prod_{j=1}^{m} W^{1-1/r,r}(\Gamma_j)$ for the same r as above. We notice at this point that the assumption $\beta_j < 2 - \frac{2}{p}$ implies that $\frac{2}{\beta_j + 2/p} > 1$, and $2 - \frac{2}{p} - \lambda_j < \beta_j$ implies $r < \frac{2}{2 - \lambda_j}$ for all j, therefore we can choose some r > 1 satisfying the assumptions of Theorem 3.4(b) and we have that $a_0 \in L^r(\Omega), f \in L^r(\Omega)$, and $u \in \prod_{j=1}^{m} W^{1-1/r,r}(\Gamma_j)$. By Theorem 3.4(b) we obtain $y \in W^{2,r}(\Omega)$ for some r > 1.

In particular, the result $y \in W^{2,r}(\Omega)$ implies $y \in L^{\infty}(\Omega)$, and hence $a_0 y \in L^p_{\vec{\beta}}(\Omega)$. We also have that $\nabla y \in W^{1,r}(\Omega)^2 \hookrightarrow L^{s_y}(\Omega)^2$ for $s_y = \frac{2r}{2-r}$ if r < 2, any $s_y < +\infty$ if r = 2 and $s_y = +\infty$ if r > 2. From this we deduce that $b \cdot \nabla y \in L^s(\Omega)$ for

$$\frac{1}{s} = \frac{1}{\hat{p}} + \frac{1}{s_y}$$

Now we use that $\vec{\beta} \geq 0$ to deduce that $b \cdot \nabla y \in L^s_{\vec{\beta}}(\Omega)$ and hence $F = -b \cdot \nabla y - a_0 y \in L^{\min\{s,p\}}_{\vec{\beta}}(\Omega)$. By applying Lemma 3.2(c), we have that $y \in W^{2,\min\{s,p\}}_{\vec{\alpha}}(\Omega)$. If $s \geq p$, the proof is complete.

Otherwise, in case $s , from Sobolev's embedding theorem, we have that <math>\nabla y \in W^{1,s}_{\vec{\beta}}(\Omega) \hookrightarrow L^{s_y}_{\vec{\beta}}(\Omega)$ for

$$\frac{1}{s_y} = \frac{1}{s} - \frac{1}{2} = \frac{s-2}{s} \iff s_y = \frac{2s}{s-2}$$

Since $\vec{\beta} \geq \vec{0}$, using that $b \in L^{\hat{p}}(\Omega)$, we have that $b \cdot \nabla y \in L^{\hat{s}}_{\vec{\beta}}$, where

$$\frac{1}{\hat{s}} = \frac{1}{s_y} + \frac{1}{\hat{p}} = \frac{1}{s} - \left(\frac{1}{2} - \frac{1}{\hat{p}}\right).$$
(3.5)

By applying Lemma 3.2(c), we have that $y \in W^{2,\min\{p,\hat{s}\}}_{\vec{\beta}}(\Omega)$. If $\hat{s} \geq 2$, the proof is complete. Otherwise, we redefine $s := \hat{s}$ and repeat the last step. Since at each iteration we subtract the positive constant $\frac{1}{2} - \frac{1}{\hat{p}}$, the proof will end in a finite number of steps.

We conjecture that the result of Theorem 3.4(c) holds for $p \in (1, \hat{p}]$, but the proof is limited to $p \leq 2$.

Notice that the operator \mathcal{A}^* is different from \mathcal{A} , and hence the results in Theorem 3.4 are not immediately applicable. For the adjoint state equation, we will need another assumption on $b \cdot n$, which is a result of the boundary term obtained due to integration by parts.

Theorem 3.5. Suppose Assumptions 2.1 and 3.1 hold. Consider $f \in H^1(\Omega)'$ and $g \in H^{1/2}(\Gamma)'$ and let $\varphi \in H^1(\Omega)$ be the unique solution of

$$\langle \mathcal{A}^* \varphi, z \rangle_{\Omega} = \langle f, z \rangle_{\Omega} + \langle g, z \rangle_{\Gamma} \quad \forall z \in H^1(\Omega).$$
(3.6)

(a) If $a_0, \nabla \cdot b \in L^q(\Omega)$, $b \cdot n \in L^{q_{\Gamma}}(\Gamma) \cap H^{t-3/2}(\Gamma)$, $f \in H^{2-t}(\Omega)'$, and $g \in \prod_{j=1}^m H^{t-3/2}(\Gamma_j)$ for $1 < t < 1 + \lambda$, $t \leq 2$, $q = \frac{2}{3-t}$, and $q_{\Gamma} = \min\{2, 1/(2-t)\}$, then $\varphi \in H^t(\Omega)$, and there exists a constant $C_{\mathcal{A}^*, t} > 0$ such that

$$\|\varphi\|_{H^t(\Omega)} \le C_{\mathcal{A}^*,t} \Big(\|f\|_{H^{2-t}(\Omega)'} + \sum_{j=1}^m \|g\|_{H^{t-3/2}(\Gamma_j)} \Big)$$

(b) If a_0 , $\nabla \cdot b$, $f \in L^r(\Omega)$, and $g, b \cdot n \in \prod_{j=1}^m W^{1-1/r,r}(\Gamma_j)$ for some $r \in (1, \hat{p}]$ satisfying $r < \frac{2}{2-\lambda}$ in case of $\lambda < 2$, then $\varphi \in W^{2,r}(\Omega)$, and there exists a constant $C_{\mathcal{A}^*,r} > 0$ such that

$$\|\varphi\|_{W^{2,r}(\Omega)} \le C_{\mathcal{A}^*,r} \Big(\|f\|_{L^r(\Omega)} + \sum_{j=1}^m \|g\|_{W^{1-1/r,r}(\Gamma_j)} \Big).$$

(c) If a_0 , $\nabla \cdot b$, $f \in L^p_{\vec{\beta}}(\Omega)$, and $b \cdot n$, $g \in \prod_{j=1}^m W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)$ for some $p \in (1,2]$ and some $\vec{\beta}$ such that $2 - \frac{2}{p} - \lambda_j < \beta_j < 2 - \frac{2}{p}$, $\beta_j \ge 0$, for all $j \in \{1, \ldots, m\}$, then $\varphi \in W^{2,p}_{\vec{\beta}}(\Omega)$ and there exists a constant $C_{\mathcal{A}^*,\vec{\beta},p} > 0$ such that

$$\|\varphi\|_{W^{2,p}_{\vec{\beta}}(\Omega)} \le C_{\mathcal{A}^*,\vec{\beta},p} \Big(\|f\|_{L^2_{\vec{\beta}}(\Omega)} + \sum_{j=1}^m \|g\|_{W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)} \Big).$$

Proof. The expression for $\langle \mathcal{A}^* \varphi, z \rangle_{\Omega}$ is derived in Lemma 2.5. Using the product rule, we have that the function φ satisfies

$$\begin{split} \int_{\Omega} \sum_{i,k=1}^{2} a_{ki} \partial_{x_{i}} \varphi \partial_{x_{k}} z \, \mathrm{d}x &- \int_{\Omega} (b \cdot \nabla \varphi) z \, \mathrm{d}x + \int_{\Omega} a_{0} \varphi z \, \mathrm{d}x \\ &= \int_{\Omega} (\nabla \cdot b) \varphi z \, \mathrm{d}x - \int_{\Gamma} \varphi (b \cdot n) z \, \, \mathrm{d}x + \langle f, z \rangle_{\Omega} + \langle g, z \rangle_{\Gamma} \end{split}$$

and we can apply Theorem 3.4 to this problem provided $\varphi \nabla \cdot b$ and $\varphi b \cdot n$ are in the appropriate spaces.

Notice that statement (a) for t = 2 is the same than statement (b) for r = 2. We will prove (a) for t < 2, and refer to (b) for t = 2.

Step 1: First, we prove $W^{1,\delta}(\Omega)$ regularity for some $\delta > 2$.

Let us write the equation as

$$\begin{cases} A\varphi + \varphi = f + \varphi \nabla \cdot b + b \cdot \nabla \varphi - a_0 \varphi + \varphi & \text{in } \Omega \\ \partial_{n_A} \varphi = -b \cdot n\varphi + g & \text{on } \Gamma. \end{cases}$$

This is a Neumann problem posed on a Lipschitz domain. We will apply the regularity results in [22]. To that end, we first investigate the existence of $r_f > 2$ and $q_{\Gamma} > 1$ such that $f \in W^{1,r'_f}(\Omega)'$ and $b \cdot n \in L^{q_{\Gamma}}(\Gamma)$. In each of the three cases, we have:

- (a) $f \in H^{2-t}(\Omega)' \hookrightarrow W^{1,r'_f}(\Omega)'$ for $r_f = \frac{2}{2-t} > 2$ since 1 < t < 2. The exponent q_{Γ} is given in the theorem.
- (b) $f \in L^r(\Omega) \hookrightarrow W^{1,r'_f}(\Omega)'$ for $r_f = \frac{2r}{2-r} > 2$ if 1 < r < 2 and all $r_f < +\infty$ if $r \ge 2$. In this case we take $q_{\Gamma} = r > 1$.
- if $r \ge 2$. In this case we take $q_{\Gamma} = r > 1$. (c) $f \in L^p_{\vec{\beta}}(\Omega) \subset L^r(\Omega)$ for $r < \frac{2}{\beta_j + \frac{2}{p}}$ for all $j \in \{1, \dots, m\}$. The condition $\beta_j < 2 - \frac{2}{p}$ implies that $\frac{2}{\beta_j + \frac{2}{p}} > 1$, so we can choose r > 1 and $L^r(\Omega) \hookrightarrow W^{1,r'_f}(\Omega)'$ for $r_f = \frac{2r}{2-r} > 2$. Therefore $f \in W^{1,r'_f}(\Omega)'$ for all $2 < r_f < \frac{2p}{2-(1-\beta_j)p}$. In this case we take $q_{\Gamma} = r > 1$.

Note that also in each of the three cases we have different assumptions on a_0 and $\nabla \cdot b$, but in any case there exists $q_0 > 1$ such that $a_0, \nabla \cdot b \in L^{q_0}(\Omega)$.

Let us check that also $F = \varphi \nabla \cdot b + b \cdot \nabla \varphi - a_0 \varphi + \varphi \in W^{1,r'_{\Omega}}(\Omega)'$ for some $r_{\Omega} > 2$. To this end define r_{φ} , s_{Ω} and r_{Ω} by

$$\frac{1}{s_{\Omega}} = \frac{1}{r_{\varphi}} = \min\left\{\frac{1}{2}\left(1 - \frac{1}{q_0}\right), \frac{1}{2} - \frac{1}{p}\right\} \in (0, \frac{1}{2}) \text{ and } \frac{1}{r_{\Omega}} = \frac{1}{2} - \frac{1}{s_{\Omega}} \in (0, \frac{1}{2})$$

such that

$$\frac{1}{r_{\varphi}} + \frac{1}{q_0} + \frac{1}{s_{\Omega}} \leq 1 \text{ and } \frac{1}{p} + \frac{1}{2} + \frac{1}{s_{\Omega}} \leq 1$$

and $W^{1,r'_{\Omega}}(\Omega) \hookrightarrow L^{s_{\Omega}}(\Omega)$. Using Lemma 2.6, we have that $\varphi \in H^{1}(\Omega) \hookrightarrow L^{r_{\varphi}}(\Omega)$. Denote $S = \{z \in W^{1,r'_{\Omega}}(\Omega) : ||z||_{W^{1,r'_{\Omega}}(\Omega)} = 1\}$. Then

$$\begin{split} \|F\|_{W^{1,r_{\Omega}'}(\Omega)'} &= \sup_{z \in S} \int_{\Omega} \left(\varphi \nabla \cdot b + b \cdot \nabla \varphi + a_{0}\varphi - \varphi\right) z \, \mathrm{d}x \\ &\leq C \sup_{z \in S} \left(\|\varphi\|_{L^{r_{\varphi}}(\Omega)} \|1 + a_{0} + \nabla \cdot b\|_{L^{q_{0}}(\Omega)} + \|b\|_{L^{p}(\Omega)} \|\nabla \varphi\|_{L^{2}(\Omega)} \right) \|z\|_{L^{s_{\Omega}}(\Omega)} \\ &\leq C_{r_{\Omega}} \sup_{z \in S} \left(\|\varphi\|_{L^{r_{\varphi}}(\Omega)} (\|1 + a_{0} + \nabla \cdot b\|_{L^{q_{0}}(\Omega)}) + \|b\|_{L^{p}(\Omega)} \|\nabla \varphi\|_{L^{2}(\Omega)} \right) \|z\|_{W^{1,r_{\Omega}'}(\Omega)} \\ &= C_{r_{\Omega}} \left(\|\varphi\|_{L^{r_{\varphi}}(\Omega)} (\|1 + a_{0} + \nabla \cdot b\|_{L^{q_{0}}(\Omega)}) + \|b\|_{L^{p}(\Omega)} \|\nabla \varphi\|_{L^{2}(\Omega)} \right). \end{split}$$

On the boundary, we want to check that $b \cdot n\varphi \in W^{-1/r_{\Gamma},r_{\Gamma}}(\Gamma) = W^{1/r_{\Gamma},r'_{\Gamma}}(\Gamma)'$ for some $r_{\Gamma} > 2$. To this end, define \hat{r}_{φ} , s_{Γ} and r_{Γ} by

$$\frac{1}{s_{\Gamma}} = \frac{1}{\hat{r}_{\varphi}} = \frac{1}{2} \left(1 - \frac{1}{q_{\Gamma}} \right) \in (0, \frac{1}{2}) \text{ and } \frac{1}{r_{\Gamma}} = \frac{1}{2} \left(1 - \frac{1}{s_{\Gamma}} \right) \in (0, \frac{1}{2})$$

such that

$$\frac{1}{\hat{r}_{\varphi}} + \frac{1}{q_{\Gamma}} + \frac{1}{s_{\Gamma}} = 1$$

and $W^{1/r_{\Gamma},r'_{\Gamma}}(\Gamma) \hookrightarrow L^{s_{\Gamma}}(\Gamma)$. From Lemma 2.3 and the trace theorem, we have that $\varphi \in H^{1/2}(\Gamma) \hookrightarrow L^{\hat{r}_{\varphi}}(\Gamma)$. Denote $S = \{z \in W^{1/r_{\Gamma},r'_{\Gamma}}(\Gamma) : \|z\|_{W^{1/r_{\Gamma},r'_{\Gamma}}(\Gamma)} = 1\}$. Then

$$\begin{aligned} \|b \cdot n\varphi\|_{W^{-1/r_{\Gamma},r_{\Gamma}}(\Gamma)} &= \sup_{z \in S} \int_{\Gamma} b \cdot n\varphi z \, \mathrm{d}x \le \sup_{z \in S} \|b \cdot n\|_{L^{q_{\Gamma}}(\Gamma)} \|\varphi\|_{L^{\hat{r}_{\varphi}}(\Gamma)} \|z\|_{L^{s_{\Gamma}}(\Gamma)} \\ &\le C_{r_{\Gamma}} \sup_{z \in S} \|b \cdot n\|_{L^{q_{\Gamma}}(\Gamma)} \|\varphi\|_{L^{\hat{r}_{\varphi}}(\Gamma)} \|z\|_{W^{1/r_{\Gamma},r'_{\Gamma}}(\Gamma)} (\Gamma) = C_{r_{\Gamma}} \|\varphi\|_{L^{\hat{r}_{\varphi}}(\Gamma)} \|b \cdot n\|_{L^{q_{\Gamma}}(\Gamma)}. \end{aligned}$$

Noting that for a general Lipschitz domain the $W^{1,\delta}(\Omega)$ regularity is limited to $\delta \leq 4$, see [22], from the previous estimates, we can deduce that, for $\delta = \min\{4, r_f, r_\Omega, r_\Gamma\} > 2$, $\varphi \in W^{1,\delta}(\Omega)$.

Step 2: Let us check that $\varphi \nabla \cdot b$ and $\varphi b \cdot n$ satisfy the regularity assumptions for the source and the Neumann data respectively of the different cases of Theorem 3.4.

(a) On one hand $\varphi \in W^{1,\delta}(\Omega) \hookrightarrow L^{\infty}(\Omega)$ and the assumption $\nabla \cdot b \in L^q(\Omega)$ imply $\varphi \nabla \cdot b \in L^q(\Omega) \hookrightarrow H^{2-t}(\Omega)'$, by the definition of q. On the boundary, by the trace theorem $\varphi \in W^{1-1/\delta,\delta}(\Gamma)$. If $1 < t \leq 3/2$, then we use that $W^{1-1/\delta,\delta}(\Gamma) \hookrightarrow L^{\infty}(\Gamma)$ to conclude that $\varphi b \cdot n \in L^{q_{\Gamma}}(\Gamma) \hookrightarrow H^{t-3/2}(\Gamma)$. The last inclusion follows by duality and the Sobolev imbedding $H^{3/2-t}(\Gamma) \hookrightarrow L^{\frac{1}{t-1}}(\Gamma)$. If 3/2 < t < 2, we use that $W^{1-1/\delta,\delta}(\Gamma) \hookrightarrow H^{s_1}(\Gamma)$ for $s_1 = 1 - 1/\delta > 1/2$. Since we are assuming that $b \cdot n \in H^{s_2}(\Gamma)$ with $s_2 = t - 3/2 \in (0, 1/2)$, from the trace theorem and the multiplication theorem [23, Theorem 7.4], we have that $\varphi b \cdot n \in H^{t-3/2}(\Gamma)$.

The result follows from Theorem 3.4(a).

(b) Using again that $\varphi \in L^{\infty}(\Omega)$, we readily deduce that $\varphi \nabla \cdot b \in L^{r}(\Omega)$. Let us check that $\varphi b \cdot n \in \prod_{i=1}^{m} W^{1-1/r,r}(\Gamma_{j})$.

For all $j \in \{1, \ldots, m\}$, by the trace theorem and the assumption on $b \cdot n$ we deduce the existence of $B_j \in W^{1,r}(\Omega)$ such that the trace of B_j on Γ_j is $b \cdot n$.

Suppose first that $r \leq 2$. Then, a straightforward application of the multiplication Lemma 3.6 below (in the case $\beta_j = 0$) yields $\varphi B_j \in W^{1,r}(\Omega)$, and hence, its trace on Γ_j belongs to $W^{1-1/r,r}(\Gamma_j)$. Therefore, $\varphi b \cdot n \in \prod_{j=1}^m W^{1-1/r,r}(\Gamma_j)$ and the result follows from Theorem 3.4(b).

If r > 2, from the previous paragraph we have that $\varphi \in W^{2,2}(\Omega) \hookrightarrow W^{1,\delta}(\Omega)$ for all $\delta < +\infty$. Repeating the previous argument, we obtain the desired result.

(c) Since $\varphi \in L^{\infty}(\Omega)$ and $\nabla \cdot b \in L^{p}_{\vec{\beta}}(\Omega)$, we have that $\varphi \nabla \cdot b \in L^{p}_{\vec{\beta}}(\Omega)$. Next, we show that $\varphi b \cdot n \in \prod_{j=1}^{m} W^{1-1/p,p}_{\vec{\beta}}(\Gamma_{j})$.

For all $j \in \{1, \ldots, m\}$, by the trace theorem and the assumption on $b \cdot n$ we deduce the existence of $B_j \in W^{1,p}_{\vec{\beta}}(\Omega)$ such that the trace of B_j on Γ_j is $b \cdot n$. Since $p \leq 2 < \delta$, a straightforward application of the multiplication Lemma 3.6 below yields $\varphi B_j \in W^{1,p}_{\vec{\beta}}(\Omega)$, and hence, its trace on Γ_j belongs to $W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)$. Therefore, $\varphi b \cdot n \in \prod_{j=1}^m W^{1-1/p,p}_{\vec{\beta}}(\Gamma_j)$ and the result follows from Theorem 3.4(c).

It remains to prove the multiplication theorem used in the proofs of cases (b) and (c) in Theorem 3.5.

Lemma 3.6. (A multiplication theorem in weighted Sobolev spaces) Let $1 < q < +\infty$. Consider $\varphi \in W^{1,\delta}(\Omega)$ for some $\delta > \max\{2,q\}$ and $\psi \in W^{1,q}_{\vec{\beta}}(\Omega)$ for some $\vec{\beta}$ such that $2 - \frac{2}{q} - \lambda_j < \beta_j < 2 - \frac{2}{q}, \beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$. Then $\psi \varphi \in W^{1,q}_{\vec{\beta}}(\Omega)$.

Proof. Since $\delta > 2$, $\varphi \in L^{\infty}(\Omega)$. Also it is clear that $\psi \in L^{q}_{\vec{\beta}}(\Omega)$, and hence $\psi \varphi \in L^{q}_{\vec{\beta}}(\Omega)$.

Let us check that also $|\nabla(\psi\varphi)| \in L^q_{\vec{\beta}}(\Omega)$. We write $\nabla(\psi\varphi) = \varphi\nabla\psi + \psi\nabla\varphi$. Using again that $\varphi \in L^{\infty}(\Omega)$ it is immediate to deduce that $|\nabla\psi| \in L^q_{\vec{\beta}}(\Omega)$ implies that $|\varphi\nabla\psi| \in L^q_{\vec{\alpha}}(\Omega)$.

Checking that the term $\psi |\nabla \varphi| \in L^q_{\vec{\beta}}(\Omega)$ is more involved. By localizing the problem at corner x_i , and applying Hölder's inequality we obtain

$$\int_{\Omega_{R_j}} (r^{\beta_j} \psi |\nabla \varphi|)^q \, \mathrm{d}x \le \|r^{\beta_j} \psi\|_{L^{\frac{q\delta}{\delta-q}}(\Omega_{R_j})}^q \|\nabla \varphi\|_{L^{\delta}(\Omega_{R_j})}^q,$$

and therefore it is sufficient to prove that $r^{\beta_j}\psi \in L^{\frac{q\delta}{\delta-q}}(\Omega_{R_j})$. Let us introduce $1 \leq q_{\delta} < q$ and $2 \leq q_{\delta}^* < +\infty$ such that

$$\frac{1}{q_{\delta}^{*}} = \min\left\{\frac{1}{2}, \frac{1}{q} - \frac{1}{\delta}\right\} \text{ and } \frac{1}{q_{\delta}} = \min\left\{1, \frac{1}{q} + \frac{1}{2} - \frac{1}{\delta}\right\} = \frac{1}{q_{\delta}^{*}} + \frac{1}{2}$$

so that $q_{\delta}^* \geq \frac{q\delta}{\delta-q}$, and $W^{1,q_{\delta}}(\Omega_{R_j}) \hookrightarrow L^{q_{\delta}^*}(\Omega_{R_j}) \hookrightarrow L^{\frac{q\delta}{\delta-q}}(\Omega_{R_j})$. We are going to prove that $r^{\beta_j} \psi \in W^{1,q_{\delta}}(\Omega_{R_j})$.

First of all we notice that $\nabla(r^{\beta_j}\psi) = r^{\beta_j}\nabla\psi + \psi\nabla r^{\beta_j}$. By definition of $W^{1,q}_{\vec{\beta}}(\Omega)$, we have that $r^{\beta_j}|\nabla\psi| \in L^q(\Omega_{R_j}) \hookrightarrow L^{q_\delta}(\Omega_{R_j})$.

For the second term we notice that $|\psi \nabla r^{\beta_j}| \sim r^{\beta_j - 1}\psi$. Since $1 - 2/q_{\delta} = \max\{-1, \frac{2}{\delta} - \frac{2}{q}\} < 0 \leq \beta_j$, we have that $W^{1,q_{\delta}}_{\vec{\beta}}(\Omega_{R_j}) \hookrightarrow L^{q_{\delta}}_{\vec{\beta}-1}(\Omega_{R_j})$; see e.g. [19, Lemma 2.29(i)]. We deduce that $\psi \in W^{1,q}_{\vec{\beta}}(\Omega_{R_j}) \hookrightarrow W^{1,q_{\delta}}_{\vec{\beta}}(\Omega_{R_j}) \hookrightarrow L^{q_{\delta}}_{\vec{\beta}-1}(\Omega_{R_j})$. This means that $r^{\beta_j - 1}\psi \in L^{q_{\delta}}(\Omega_{R_j})$, and we gather that $|\psi \nabla r^{\beta_j}| \in L^{q_{\delta}}(\Omega_{R_j})$.

Therefore $\nabla(r^{\beta_j}\psi) \in L^{q_\delta}(\Omega_{R_j})$, so we have that $r^{\beta_j}\psi \in W^{1,q_\delta}(\Omega_{R_j})$.

Using this, we conclude that $\psi |\nabla \varphi| \in L^q_{\vec{\beta}}(\Omega)$ and consequently $|\nabla(\psi \varphi)| \in L^q_{\vec{\beta}}(\Omega)$, which leads to the desired result.

4. Discretization

Consider a family of regular triangulations $\{\mathcal{T}_h\}$ graded with mesh grading parameters $\mu_j \in (0, 1], j \in \{1, \ldots, m\}$ in the sense of [24, Section 3.1], see also [25]. As usual, $Y_h \subset H^1(\Omega) \cap C(\overline{\Omega})$ is the space of continuous piecewise linear functions.

Lemma 4.1. There exists a constant $c_{\vec{\mu}} > 0$ such that

$$\|\psi - I_h \psi\|_{H^1(\Omega)} \le c_{\vec{\mu}} h^s \|\psi\|_{W^{2,2}_{\vec{a}}(\Omega)} \ \forall \psi \in W^{2,2}_{\vec{\beta}}(\Omega),$$

where I_h is the Lagrange interpolation operator, the vector $\vec{\beta}$ satisfies that $1 - \lambda_j < \beta_j < 1$ and $\beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$ and the exponent s satisfies that $s \le 1$ and $s < \frac{\lambda_j}{\mu_i}$ for all $j \in \{1, \ldots, m\}$.

Proof. The case $\mu_j = 1$ (quasi-uniform mesh) is classical. For $\mu_j < \lambda_j$ see [4, Lemma 4.1]. The case $\lambda_j \leq \mu_j < 1$ can be proved with the same techniques and the additional idea that $h_T \sim h^s r_T^{1-s\mu_j}$, $1-s\mu_j = \beta_j > 1-\lambda_j$; see equation (3.14) in [24, Theorem 3.2], where it was used for a Dirichlet problem. \Box

Define the bilinear form $a(y,z) = \langle Ay, z \rangle_{\Omega}$. For a datum $u \in H^{1/2}(\Gamma)'$, the discrete state equation reads

$$a(y_h, z_h) = \langle u, z_h \rangle_{\Gamma} \ \forall z_h \in Y_h.$$

$$(4.1)$$

Existence and uniqueness of the solution of this equation is not immediate since $a(\cdot, \cdot)$ is not coercive over Y_h .

Theorem 4.2. There exists $h_0 > 0$ that depends on A, b, a_0 , Ω and the mesh grading parameter $\vec{\mu}$, such that the system (4.1) has a unique solution for every $h \leq h_0$ and every $u \in H^{1/2}(\Gamma)'$. Further, there exists a constant K_0 that depends on A, b, a_0 , Ω and is independent of $\vec{\mu}$ and h such that

$$\|y_h\|_{H^1(\Omega)} \le K_0 \|\mathcal{A}^{-1}u\|_{H^1(\Omega)} \ \forall h \le h_0.$$
(4.2)

The scheme of the proof is similar to that of [9, Lemma 3.1] for distributed control problems with homogeneous Dirichlet boundary conditions, but that proof is done for quasi-uniform meshes and uses this fact explicitly; see equations (3.8) and (3.9) in [9]. Since the mesh grading and the boundary terms imply some extra technicalities, we include a complete proof for the convenience of the reader.

Proof. Due to the linearity of the system, to show existence it is sufficient to prove uniqueness of solution in the case u = 0. Suppose $y_h \in Y_h$ satisfies

$$a(y_h, z_h) = 0 \ \forall z_h \in Y_h.$$

$$(4.3)$$

Taking $z_h = y_h$ and using Gårding's inequality established in Lemma 2.3, we have that

$$0 = a(y_h, y_h) = \langle \mathcal{A}y_h, y_h \rangle_{\Omega} \ge \frac{\Lambda}{8C_E^2} \|y_h\|_{H^1(\Omega)}^2 - C_{\Lambda, E, b} \|y_h\|_{L^2(\Omega)}^2.$$

Therefore

$$\|y_h\|_{H^1(\Omega)} \le 2C_E \sqrt{\frac{2C_{\Lambda, E, b}}{\Lambda}} \|y_h\|_{L^2(\Omega)}.$$
(4.4)

Since $y_h \in L^2(\Omega) \subset L^2_{\vec{\beta}}(\Omega)$ for all $\vec{\beta} \geq \vec{0}$ such that $1 - \lambda_j < \beta_j$ for all $j \in \{1, \ldots, m\}$, from Theorem 3.5(c), we have that there exists a unique $\psi \in W^{2,2}_{\vec{\beta}}(\Omega)$ such that

$$a(z,\psi) = \int_{\Omega} y_h z \mathrm{d}x \ \forall z \in H^1(\Omega)$$
(4.5)

and there exists a constant $C_{\mathcal{A}^*,\vec{\beta}}$ such that

$$\|\psi\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \le C_{\mathcal{A}^*,\vec{\beta}} \|y_h\|_{L^2_{\vec{\beta}}(\Omega)}.$$

Let us denote $\hat{\psi}_h \in Y_h$ the Ritz–Galerkin projection of ψ onto Y_h in the sense of $H^1(\Omega)$, i.e., $\hat{\psi}_h$ is the unique solution of

$$\int_{\Omega} (\nabla \hat{\psi}_h \nabla z_h + \hat{\psi}_h z_h) \mathrm{d}x = \int_{\Omega} (\nabla \psi \nabla z_h + \psi z_h) \mathrm{d}x \ \forall z_h \in Y_h.$$

From [5, Eq. (4.2)], Theorem 3.5(c), and the embedding $L^2(\Omega) \hookrightarrow L^2_{\vec{\beta}}(\Omega)$, with embedding constant 1 due to the choice $R_j \leq 1$, we have that there exists a constant $\hat{c}_{\vec{\mu}}$ such that

$$\|\psi - \hat{\psi}_h\|_{H^1(\Omega)} \le \hat{c}_{\vec{\mu}} h^s \|\psi\|_{W^{2,2}_{\beta}(\Omega)} \le \hat{c}_{\vec{\mu}} C_{\mathcal{A}^*,\beta} h^s \|y_h\|_{L^2_{\vec{\beta}}(\Omega)} \le \hat{c}_{\vec{\mu}} C_{\mathcal{A}^*,\beta} h^s \|y_h\|_{L^2(\Omega)},$$
(4.6)

where $s \leq 1$ and $s < \frac{\lambda_j}{\mu_j}$ for all $j \in \{1, \ldots, m\}$; see Lemma 4.1. Taking $z = y_h$ in the adjoint Eq. (4.5), and $z_h = \hat{\psi}_h$ in the homogeneous discrete Eq. (4.3), we deduce

$$\begin{aligned} \|y_h\|_{L^2(\Omega)}^2 &= a(y_h, \psi) = a(y_h, \psi - \hat{\psi}_h) \le \|\mathcal{A}\| \|y_h\|_{H^1(\Omega)} \|\psi - \hat{\psi}_h\|_{H^1(\Omega)} \\ &\le \hat{c}_{\vec{\mu}} C_{\mathcal{A}^*, \vec{\beta}} \|\mathcal{A}\| \|y_h\|_{H^1(\Omega)} h^s \|y_h\|_{L^2(\Omega)}. \end{aligned}$$

Along the proof we will denote $\|\mathcal{A}\| = \|\mathcal{A}\|_{\mathcal{L}(H^1(\Omega), H^1(\Omega)')}$. Choosing h_0 such that

$$\hat{c}_{\vec{\mu}}C_{\mathcal{A}^*,\vec{\beta}} \|\mathcal{A}\| h_0^s = \frac{1}{2} \frac{1}{2C_E} \sqrt{\frac{\Lambda}{2C_{\Lambda,E,b}}},\tag{4.7}$$

we have that, for all $h \leq h_0$

$$\|y_h\|_{L^2(\Omega)} \le \frac{1}{2} \frac{1}{2C_E} \sqrt{\frac{\Lambda}{2C_{\Lambda,E,b}}} \|y_h\|_{H^1(\Omega)}.$$

Using this and estimate (4.4), we deduce that

$$||y_h||_{H^1(\Omega)} \le \frac{1}{2} ||y_h||_{H^1(\Omega)} \ \forall h \le h_0,$$

and hence $y_h = 0$.

Take now $u \in H^{1/2}(\Gamma)'$ and denote $y = \mathcal{A}^{-1}u$. For $h \leq h_0$, let y_h be the solution of (4.1). Taking $z = y_h$ in the adjoint Eq. (4.5), and $z_h = \hat{\psi}_h$ in the discrete Eq. (4.1), we deduce

$$\begin{split} \|y_h\|_{L^2(\Omega)}^2 &= a(y_h, \psi) = a(y_h, \psi - \hat{\psi}_h) + \langle u, \hat{\psi}_h \rangle_{\Gamma} = a(y_h, \psi - \hat{\psi}_h) + a(y, \hat{\psi}_h) \\ &\leq \|\mathcal{A}\| \left(\|y_h\|_{H^1(\Omega)} \|\psi - \hat{\psi}_h\|_{H^1(\Omega)} + \|y\|_{H^1(\Omega)} \|\hat{\psi}_h\|_{H^1(\Omega)} \right) \\ &\leq \hat{c}_{\vec{\mu}} C_{\mathcal{A}^*, \vec{\beta}} \|\mathcal{A}\| \|y_h\|_{H^1(\Omega)} h^s \|y_h\|_{L^2(\Omega)} + c_{\vec{\beta}} C_{\mathcal{A}^*, \vec{\beta}} \|\mathcal{A}\| \|y\|_{H^1(\Omega)} \|y_h\|_{L^2(\Omega)}, \end{split}$$

where we have used that $\|\hat{\psi}_h\|_{H^1(\Omega)} \leq \|\psi\|_{H^1(\Omega)} \leq \hat{c}_{\vec{\beta}} \|\psi\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \leq c_{\vec{\beta}} C_{\mathcal{A}^*,\vec{\beta}} \|y_h\|_{L^2(\Omega)}$; see [19, Lemma 2.29(i)] for the embedding $W^{2,2}_{\vec{\beta}}(\Omega) \hookrightarrow H^1(\Omega)$. Now, using that $h \leq h_0$ and (4.7) we have

$$\|y_h\|_{L^2(\Omega)} \le \frac{1}{2} \frac{1}{2C_E} \sqrt{\frac{\Lambda}{2C_{\Lambda,E,b}}} \|y_h\|_{H^1(\Omega)} + c_{\vec{\beta}} C_{\mathcal{A}^*,\vec{\beta}} \|\mathcal{A}\| \|y\|_{H^1(\Omega)}$$

and applying Young's inequality we deduce

$$\|y_h\|_{L^2(\Omega)}^2 \le \frac{1}{16} \frac{1}{C_E^2} \frac{\Lambda}{C_{\Lambda,E,b}} \|y_h\|_{H^1(\Omega)}^2 + 2c_{\vec{\beta}}^2 C_{\mathcal{A}^*,\vec{\beta}}^2 \|\mathcal{A}\|^2 \|y\|_{H^1(\Omega)}^2.$$
(4.8)

Using Gårding's inequality, the discrete Eq. (4.1) and $y = \mathcal{A}^{-1}u$, we infer

$$\frac{\Lambda}{8C_E^2} \|y_h\|_{H^1(\Omega)}^2 - C_{\Lambda,E,b} \|y_h\|_{L^2(\Omega)}^2 \le a(y_h, y_h)
= \langle u, y_h \rangle_{\Gamma} = a(y, y_h) \le \|\mathcal{A}\| \|y\|_{H^1(\Omega)} \|y_h\|_{H^1(\Omega)}.$$
(4.9)

Multiplying (4.8) by $C_{\Lambda,E,b}$ and using the resulting inequality in (4.9), we obtain

$$\begin{aligned} \frac{\Lambda}{16C_E^2} \|y_h\|_{H^1(\Omega)}^2 &\leq 2c_{\vec{\beta}}^2 C_{\mathcal{A}^*,\vec{\beta}}^2 \|\mathcal{A}\|^2 \|y\|_{H^1(\Omega)}^2 + \|\mathcal{A}\| \|y\|_{H^1(\Omega)} \|y_h\|_{H^1(\Omega)} \\ &\leq 2c_{\vec{\beta}}^2 C_{\mathcal{A}^*,\vec{\beta}}^2 \|\mathcal{A}\|^2 \|y\|_{H^1(\Omega)}^2 + \frac{8C_E^2}{\Lambda} \|\mathcal{A}\|^2 \|y\|_{H^1(\Omega)}^2 + \frac{\Lambda}{32C_E^2} \|y_h\|_{H^1(\Omega)}^2, \end{aligned}$$

where in the second step we have used Young's inequality. Gathering the terms with $\|y_h\|_{H^1(\Omega)}^2$ and taking the square root, we finally obtain:

$$\|y_h\|_{H^1(\Omega)} \le \frac{C_E}{4\sqrt{2\Lambda}} \|\mathcal{A}\| \left(2c_{\vec{\beta}}^2 C_{\mathcal{A}^*,\vec{\beta}}^2 + \frac{8C_E^2}{\Lambda} \right)^{1/2} \|\mathcal{A}^{-1}u\|_{H^1(\Omega)}$$

Notice that the constant depends on $\vec{\beta}$, which is itself limited by the value of $\vec{\lambda}$, and hence the constant will finally depend on $\vec{\lambda}$.

Theorem 4.3. There exists $h_0^* > 0$ that depends on A, b, a_0 , Ω and the mesh grading parameter $\vec{\mu}$, such that the discrete adjoint problem

$$a(z_h, \varphi_h) = \langle y, z_h \rangle_\Omega \ \forall z_h \in Y_h \tag{4.10}$$

has a unique solution for every $y \in H^1(\Omega)'$ and every $0 < h \leq h_0^*$. Further, there exists a constant K_0^* that depends on A, b, a_0 , Ω and is independent of $\vec{\mu}$ and h such that

$$\|\varphi_h\|_{H^1(\Omega)} \le K_0^* \|(\mathcal{A}^*)^{-1}y\|_{H^1(\Omega)} \ \forall h < h_0^*.$$
(4.11)

Proof. Existence and uniqueness of solution of the discrete adjoint Eq. (4.10) follows for all $0 < h < h_0$ due to the finite-dimensional character of the problem. To get the estimate (4.11), we follow the steps of the proof of Theorem 4.2. Notice that in this case, the value of h_0^* , which is used explicitly in the proof, may be different from the value of h_0 provided in (4.7).

The following estimate is an immediate consequence of the previous results, Lemma 2.6, Corollary 2.7 and the trace theorem.

Corollary 4.4. Let $\bar{h} = \min\{h_0, h_0^*\}$ with h_0 from Theorem 4.2 and h_0^* from Theorem 4.3. For $u \in L^2(\Gamma)$ let $y_h \in Y_h$ be the unique solution of (4.1). There exists a constant $c_2 > 0$ that depends on the data of the problem, but not on the mesh grading parameters $\vec{\mu}$ or on h, such that, for all $h < \bar{h}$

$$\|y_h\|_{L^2(\Gamma)} \le c_2 \|u\|_{L^2(\Gamma)}.$$
(4.12)

Proof. Let us denote C_{TR} the norm of the trace operator from $H^1(\Omega)$ to $L^2(\Gamma)$. We use Theorem 4.2, Lemma 2.6, and the fact that u can be seen as an element of $H^1(\Omega)'$ and $\|u\|_{H^1(\Omega)'} \leq C_{\text{TR}} \|u\|_{L^2(\Gamma)}$, cf. (2.4) and (2.5). A straightforward estimation shows that

$$\|y_h\|_{L^2(\Gamma)} \le C_{\mathrm{TR}} \|y_h\|_{H^1(\Omega)} \le C_{\mathrm{TR}} K_0 \|\mathcal{A}^{-1}u\|_{H^1(\Omega)} \le C_{\mathrm{TR}} K_0 \|\mathcal{A}^{-1}\|_{C_{\mathrm{TR}}} \|u\|_{L^2(\Gamma)},$$

where $\|\mathcal{A}^{-1}\|$ denotes the norm in $\mathcal{L}(H^1(\Omega)', H^1(\Omega))$. The result follows for $c_2 = C_{\mathrm{TR}}^2 K_0 \|\mathcal{A}^{-1}\|$.

Theorem 4.5 (Error estimates in the domain). For $0 < h < \overline{h}$, where \overline{h} is defined in Corollary 4.4, and $u \in H^{1/2}(\Gamma)'$, let $y_h \in Y_h$ be the solution of (4.1) and $y \in H^1(\Omega)$ be the solution of (3.1) for f = 0. There exists C > 0 that depends on A, b, a_0 , Ω but is independent of h such that

$$\|y - y_h\|_{L^2(\Omega)} \le Ch^s \|u\|_{H^{1/2}(\Gamma)'}.$$
(4.13)

If further $u \in \prod_{j=1}^{m} W_{\vec{\beta}}^{1/2,2}(\Gamma_j)$, where $1 - \lambda_j < \beta_j < 1$ and $\beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$, there exists C > 0 that depends on A, b, a_0 , Ω , $\vec{\beta}$, and the mesh grading parameter $\vec{\mu}$, but is independent of h and u such that

$$\|y - y_h\|_{L^2(\Omega)} + h^s \|y - y_h\|_{H^1(\Omega)} \le Ch^{2s} \|y\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \le Ch^{2s} \sum_{j=1}^m \|u\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)}$$

$$(4.14)$$

for all $s \leq 1$ and $s < \frac{\lambda_j}{\mu_j}$ for all $j \in \{1, \ldots, m\}$.

Furthermore, for all $f \in L^2_{\vec{\beta}}(\Omega)$ and $g \in \prod_{j=1}^m W^{1/2,2}_{\vec{\beta}}(\Gamma)$, let $\varphi \in W^{2,2}_{\vec{\beta}}(\Omega)$ be the solution of (3.6) and φ_h be the unique solution of

$$a(z_h,\varphi_h) = \int_{\Omega} f z_h \, \mathrm{d}x + \int_{\Gamma} g z_h \, \mathrm{d}x \, \forall z_h \in Y_h.$$

Then

$$\begin{aligned} \|\varphi - \varphi_h\|_{L^2(\Omega)} + h^s \|\varphi - \varphi_h\|_{H^1(\Omega)} &\leq Ch^{2s} \|\varphi\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \\ &\leq Ch^{2s} \left(\|f\|_{L^2_{\vec{\beta}}(\Omega)} + \sum_{j=1}^m \|g\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)} \right). \end{aligned}$$
(4.15)

Proof. We will prove (4.13) and (4.14). The proof of (4.15) follows the same lines.

We first prove that

$$\|y - y_h\|_{L^2(\Omega)} \le C_{\mathcal{A}^*,\beta} \hat{c}_{\vec{\mu}} \|\mathcal{A}\| h^s \|y - y_h\|_{H^1(\Omega)}$$
(4.16)

Consider $\psi \in W^{2,2}_{\beta}(\Omega)$ the solution of the adjoint problem

$$a(z,\psi) = \int_{\Omega} (y-y_h) z \mathrm{d}x \ \forall z \in H^1(\Omega)$$

and let $\psi_h \in Y_h$ be its Ritz–Galerkin projection onto Y_h in the sense of $H^1(\Omega)$, as in the proof of Theorem 4.2. We have, with (4.6), that

$$\begin{aligned} \|y - y_h\|_{L^2(\Omega)}^2 &= a(y - y_h, \psi) = a(y - y_h, \psi - \dot{\psi}_h) \\ &\leq \|\mathcal{A}\| \|y - y_h\|_{H^1(\Omega)} \|\psi - \dot{\psi}_h\|_{H^1(\Omega)} \\ &\leq C_{\mathcal{A}^*, \beta} \hat{c}_{\vec{\mu}} \|\mathcal{A}\|_{h^s} \|y - y_h\|_{H^1(\Omega)} \|y - y_h\|_{L^2(\Omega)} \end{aligned}$$

and (4.16) follows. Estimate (4.13) follows from this, Theorem 4.2 and Lemma 2.6.

Using Gårding's inequality established in Lemma 2.3, estimate (4.16), and the definition of $h_0 > 0$ in (4.7), we have that for all $h < h_0$

$$\begin{split} \frac{\Lambda}{8C_E^2} \|y - y_h\|_{H^1(\Omega)}^2 &\leq a(y - y_h, y - y_h) + C_{\Lambda, E, b} \|y - y_h\|_{L^2(\Omega)}^2 \\ &\leq a(y - y_h, y - y_h) + C_{\Lambda, E, b} \Big(C_{\mathcal{A}^*, \beta} \hat{c}_{\vec{\mu}} \|\mathcal{A}\| h^s \Big)^2 \|y - y_h\|_{H^1(\Omega)}^2 \\ &\leq a(y - y_h, y - y_h) + \frac{1}{4} \frac{\Lambda}{8C_E^2} \|y - y_h\|_{H^1(\Omega)}^2, \end{split}$$

and hence

$$\frac{3\Lambda}{32C_E^2} \|y - y_h\|_{H^1(\Omega)}^2 \le a(y - y_h, y - y_h)$$
(4.17)

Using Theorem 3.4(c) and Lemma 4.1

$$\|y - I_h y_h\|_{H^1(\Omega)} \le \hat{c}_{\vec{\mu}} h^s \|y\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \le \hat{c}_{\vec{\mu}} C_{\mathcal{A},\vec{\beta}} h^s \sum_{j=1}^m \|u\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)}.$$
 (4.18)

Using that $a(y, I_h y_h) = a(y_h, I_h y_h)$, (4.17) and the above inequality, we have that

$$\frac{3\Lambda}{32C_E^2} \|y - y_h\|_{H^1(\Omega)}^2 \le a(y - y_h, y - I_h y_h) \le \|\mathcal{A}\| \|y - y_h\|_{H^1(\Omega)} \|y - I_h y_h\|_{H^1(\Omega)} \le \hat{c}_{\vec{\mu}} C_{\mathcal{A},\vec{\beta}} \|\mathcal{A}\| h^s \|y - y_h\|_{H^1(\Omega)},$$

and the result follows.

Corollary 4.6. There exists C > 0 that depends on A, b, a_0 , Ω , and the mesh grading parameter $\vec{\mu}$, but is independent of h such that for $0 < h < \bar{h}$

$$\|y - y_h\|_{L^2(\Omega)} \le Ch^{3s/2} \|u\|_{L^2(\Gamma)} \ \forall u \in L^2(\Gamma)$$
(4.19)

for all $s \leq 1$ and $s < \frac{\lambda_j}{\mu_j}$ for all $j \in \{1, \ldots, m\}$.

Further, for all $f \in L^2_{\vec{\beta}}(\Omega)$ and $g \in \prod_{j=1}^m W^{1/2,2}_{\vec{\beta}}(\Gamma)$ and all $\theta \in (0,1)$, then

$$\|\varphi - \varphi_h\|_{H^{\theta}(\Omega)} \le Ch^{(2-\theta)s} \left(\|f\|_{L^2_{\tilde{\beta}}(\Omega)} + \sum_{j=1}^m \|g\|_{W^{1/2,2}_{\tilde{\beta}}(\Gamma_j)} \right), \quad (4.20)$$

where C is independent of θ .

Proof. If $u \in H^{1/2}(\Gamma)$, then, by (4.14) and the embedding $H^{1/2}(\Gamma) \hookrightarrow W^{1/2,2}_{\vec{\beta}}(\Gamma) \hookrightarrow \prod_{j=1}^{m} W^{1/2,2}_{\vec{\beta}}(\Gamma_j)$ for some $\vec{\beta}$ with $\beta_j \ge 0, 1 - \lambda_j < \beta_j < 1$, we obtain

$$||y - y_h||_{L^2(\Omega)} \le Ch^{2s} ||u||_{H^{1/2}(\Gamma)}.$$

The first result follows by complex interpolation between this estimate and (4.13).

The second one follows by interpolation between the estimates for $\theta = 0$ and $\theta = 1$ that follow from (4.15).

5. Analysis of the Control Problem

Now, we turn to the analysis of the control problem

(P)
$$\min_{u \in U_{\text{ad}}} J(u) := \frac{1}{2} \int_{\Omega} (y_u(x) - y_d(x))^2 \, \mathrm{d}x + \frac{\nu}{2} \int_{\Gamma} u^2(x) \, \mathrm{d}x + \int_{\Gamma} y_u(x) g_{\varphi}(x) \, \mathrm{d}x,$$

where $y_u \in H^1(\Omega)$ solves (2.7). For every $u \in H^{1/2}(\Gamma)'$, we define $\varphi_u \in H^1(\Omega)$ as the unique solution of

$$\langle z, \mathcal{A}^* \varphi_u \rangle_{\Omega} = \int_{\Omega} (y_u - y_d) z \, \mathrm{d}x + \int_{\Gamma} g_{\varphi} z \, \mathrm{d}x \forall z \in H^1(\Omega).$$

We have that

$$J'(u)v = \int_{\Omega} (\varphi_u + \nu u)v \,\mathrm{d}x.$$

 \square

Theorem 5.1. For any $y_d \in L^2(\Omega)$ and $g_{\varphi} \in L^2(\Gamma)$, problem (P) has a unique solution $\bar{u} \in U_{ad}$ and there exist $\bar{y}, \bar{\varphi} \in H^1(\Omega)$ such that

$$\begin{split} \langle \mathcal{A}\bar{y}, z \rangle_{\Omega} &= \int_{\Gamma} \bar{u}z \, \mathrm{d}x & \forall z \in H^{1}(\Omega), \\ \langle z, \mathcal{A}^{*}\bar{\varphi} \rangle_{\Omega} &= \int_{\Omega} (\bar{y} - y_{d})z \, \mathrm{d}x + \int_{\Gamma} g_{\varphi}z \, \mathrm{d}x & \forall z \in H^{1}(\Omega), \\ \int_{\Gamma} (\bar{\varphi} + \nu \bar{u})(u - \bar{u}) \, \mathrm{d}x \geq 0 & \forall u \in U_{\mathrm{ad}}, \end{split}$$

and $\bar{u} \in H^{1/2}(\Gamma)$.

If, further, $g_{\varphi} \in \prod_{j=1}^{m} W_{\vec{\beta}}^{1/2,2}(\Gamma_{j})$ for some $\vec{\beta}$ such that $1 - \lambda_{j} < \beta_{j} < 1$ and $\beta_{j} \geq 0$ for all $j \in \{1, \ldots, m\}$, then $\bar{y}, \bar{\varphi} \in W_{\vec{\beta}}^{2,2}(\Omega) \cap C(\bar{\Omega}), \bar{\varphi} \in W_{\vec{\beta}}^{3/2,2}(\Gamma) \cap C(\Gamma), \bar{u} \in C(\Gamma).$

If, moreover, the weights also satisfy $\beta_j < 1/2$, for all $j \in \{1, \ldots, m\}$ then $\overline{\varphi}, \overline{u} \in H^1(\Gamma)$.

Proof. The existence of the solution follows from the appropriate continuity properties of the involved operators that are deduced from Lemma 2.6. Uniqueness is deduced from the strict convexity of the functional. The first order optimality conditions are deduced, hence, in a standard way from the Euler-Lagrange equation $J'(\bar{u})(u-\bar{u}) \geq 0$ for all $u \in U_{\rm ad}$ and Corollary 2.7. The $H^1(\Omega)$ regularity of \bar{y} follows from Lemma 2.3 and the regularity of the adjoint state from Lemma 2.6. By the trace theorem, we have that $\varphi \in H^{1/2}(\Gamma)$. This and the projection formula

$$\bar{u}(x) = \operatorname{Proj}_{[u_a, u_b]}\left(-\frac{\bar{\varphi}(x)}{\nu}\right),\tag{5.1}$$

which follows in a standard way from the third optimality condition, imply the regularity of \bar{u} .

Suppose now that g_{φ} belongs to $L^2(\Gamma) \cap \prod_{j=1}^m W_{\vec{\beta}}^{1/2,2}(\Gamma_j)$ for some $\vec{\beta}$ such that $1 - \lambda_j < \beta_j < 1$ and $\beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$. The $W_{\vec{\beta}}^{2,2}(\Omega)$ regularity of the state and adjoint state follow from a bootstrapping argument: since $\bar{y} \in H^1(\Omega)$ and $\beta_j \ge 0$ for all j, we have that $\bar{y} - y_d \in L^2(\Omega) \hookrightarrow L^2_{\vec{\beta}}(\Omega)$. From Theorem 3.5(c) we deduce that $\bar{\varphi} \in W^{2,2}_{\vec{\beta}}(\Omega)$. This readily implies that $\bar{\varphi} \in W^{3/2,2}_{\vec{\beta}}(\Gamma)$. Using that $L^2_{\vec{\beta}}(\Omega) \subset L^r(\Omega)$ for all $1 < r < 2/(1 + \beta_j)$, we deduce from Theorem 3.5(b) that $\bar{\varphi} \in W^{2,r}(\Omega) \hookrightarrow C(\bar{\Omega})$, so $\bar{\varphi} \in C(\Gamma)$. Again the projection formula leads to $\bar{u} \in C(\Gamma)$.

If $\beta_j < 1/2$, then $2/(1 + \beta_j) > 4/3$, so there exists r > 4/3 such that $\bar{\varphi} \in W^{2,r}(\Omega) \hookrightarrow H^{3-2/r}(\Omega)$. Since 3-2/r > 3/2, by the trace theorem we have that $\bar{\varphi} \in C(\Gamma) \cap_{j=1}^m H^1(\Gamma_j) = H^1(\Gamma)$. This last equality follows because Γ is one-dimensional and polygonal. This regularity is preserved by the projection formula, and therefore $\bar{u} \in H^1(\Gamma)$.

Notice that for any polygonal domain $\lambda_j > 1/2$ for all $j \in \{1, \ldots, m\}$, so the condition $\beta_j < 1/2$ may be a constraint in the regularity of the datum g_{φ} , but it is not a constraint on the domain. Although some of the intermediate results below can be proved for $g_{\varphi} \in L^2(\Gamma)$, since the main result requires $H^1(\Gamma)$ regularity of the optimal control, in the rest of the work we will do the following assumption.

Assumption 5.2. We assume that $g_{\varphi} \in \prod_{j=1}^{m} W_{\vec{\beta}}^{1/2,2}(\Gamma_j)$ for some $\vec{\beta}$ such that $1 - \lambda_j < \beta_j < 1/2, \ \beta_j \ge 0$ for all $j \in \{1, \ldots, m\}$. We denote

$$M_d = \|y_d\|_{L^2(\Omega)} + \sum_{j=1}^m \|g_{\varphi}\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)} + 1.$$

For every $u \in L^2(\Gamma)$, we will denote $y_h(u)$ the solution of the discrete state Eq. (4.1) and $\varphi_h(u)$ the solution of

$$a(z_h, \varphi_h) = \int_{\Omega} (y_h(u) - y_d) z_h \, \mathrm{d}x \int_{\Gamma} g_{\varphi} z_h \, \mathrm{d}x \, \forall z_h \in Y_h.$$

Our discrete functional reads like

$$J_{h}(u) = \frac{1}{2} \int_{\Omega} (y_{h}(u) - y_{d})^{2} \, \mathrm{d}x + \frac{\nu}{2} \int_{\Gamma} u^{2} \, \mathrm{d}x + \int_{\Gamma} y_{h}(u) g_{\varphi} \, \mathrm{d}x$$

To discretize the control, we notice that every triangulation \mathcal{T}_h of Ω defines a segmentation \mathcal{E}_h of Γ and define $U_{h,\mathrm{ad}} = U_h \cap U_{\mathrm{ad}}$, where

$$U_h = \{ u_h \in L^2(\Gamma) : u_{h|E} \in \mathcal{P}^0(E) \; \forall E \in \mathcal{E}_h \}.$$

Here and elsewhere $\mathcal{P}^i(K)$ is the set of polynomials of degree *i* in the set *K*. For every $u \in L^1(\Gamma)$, we define $Q_h u \in U_h$ by

$$Q_h u(x) = \frac{1}{h_E} \int_E u \, \mathrm{d}x \text{ if } x \in E,$$

where $E \in \mathcal{E}_h$ and h_E is the length of E. Notice that $u \in U_{ad}$ implies $Q_h u \in U_{h,ad}$.

Lemma 5.3. For every $u \in H^1(\Gamma)$ there exists a constant C > 0 independent of h such that

$$||u - Q_h u||_{(H^1(\Gamma))'} + h||u - Q_h u||_{L^2(\Gamma)} \le Ch^2 ||u||_{H^1(\Gamma)}.$$

If Assumption 5.2 holds, then we also have that

$$\left| \int_{\Gamma} (\varphi_u + \nu u) (u - Q_h u) \, \mathrm{d}x \right| \le Ch^2 \left(\|u\|_{H^1(\Gamma)}^2 + M_d^2 \right)$$

Proof. It is well known that for every $E \in \mathcal{E}_h$ we have $||u - Q_h u||_{L^2(E)} \leq Ch_E ||u||_{H^1(E)}$. Using that $h_E \leq ch$, we have

$$\|u - Q_h u\|_{L^2(\Gamma)}^2 = \sum_{E \in \mathcal{E}_h} \|u - Q_h u\|_{L^2(E)}^2 \le C \sum_{E \in \mathcal{E}_h} h_E^2 \|u\|_{H^1(E)}^2 \le C h^2 \|u\|_{H^1(\Gamma)}^2.$$

The estimate for the norm in $H^1(\Gamma)'$ follows now by duality since $\int_{\Gamma} (u - Q_h u) w_h dx = 0$ for all $w_h \in U_h$. This estimate implies the third one taking into account that, using the same arguments as in the proof of Theorem 5.1, $\varphi_u \in H^1(\Gamma)$, and

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$$\begin{aligned} \|\varphi_u\|_{H^1(\Gamma)} &\leq C \|\varphi\|_{W^{2,2}_{\vec{\beta}}(\Omega)} \leq C \left(\|y_u - y_d\|_{L^2_{\vec{\beta}}(\Omega)} + \sum_{j=1}^m \|g_\varphi\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)} \right) \\ &\leq C \left(\|y_u\|_{L^2(\Omega)} + \|y_d\|_{L^2(\Omega)} + \sum_{j=1}^m \|g_\varphi\|_{W^{1/2,2}_{\vec{\beta}}(\Gamma_j)} \right) \leq C \left(\|u\|_{L^2(\Gamma)} + M_d \right). \end{aligned}$$

Therefore, we obtain

$$\left| \int_{\Gamma} (\varphi_u + \nu u)(u - Q_h u) \, \mathrm{d}x \right| \leq \|\varphi_u + \nu u\|_{H^1(\Gamma)} \|u - Q_h u\|_{H^1(\Gamma)'} \\ \leq C \left(M_d + \|u\|_{L^2(\Gamma)} + \nu \|u\|_{H^1(\Gamma)} \right) h^2 \|u\|_{H^1(\Gamma)}$$

and the result follows using Young's inequality.

Our discrete problems reads like

$$(\mathbf{P}_h) \min_{u_h \in U_{h,\mathrm{ad}}} J_h(u_h).$$

Existence and uniqueness of solution of problem (P_h) , as well as first order optimality conditions follow in a standard way. We state the result in the next theorem for further reference.

Theorem 5.4. For every $0 < h < \bar{h}$, problem (P_h) has a unique solution $\bar{u}_h \in U_{h,ad}$. Further, if we denote $\bar{y}_h = y_h(\bar{u}_h)$ and $\bar{\varphi}_h = \varphi_h(\bar{u}_h)$, then

$$\int_{\Gamma} (\bar{\varphi}_h + \nu \bar{u}_h) (u_h - \bar{u}_h) \, \mathrm{d}x \ge 0 \,\,\forall u_h \in U_{h,\mathrm{ad}}.$$
(5.2)

Before stating and proving the main theorem of this section, we prove two auxiliary results.

Lemma 5.5. There exists C > 0 independent of h, y_d and g_{φ} such that for all $0 < h < \overline{h}$,

$$\|\bar{y}\|_{H^{1}(\Omega)} + \|\bar{u}\|_{H^{1/2}(\Gamma)} + \|\bar{y}_{h}\|_{H^{1}(\Omega)} + \|\bar{u}_{h}\|_{L^{2}(\Gamma)} \le C\left(\|y_{d}\|_{L^{2}(\Omega)} + \|g_{\varphi}\|_{L^{2}(\Gamma)} + 1\right).$$

If, moreover, Assumption 5.2 holds, then

$$\|\bar{u}\|_{H^1(\Gamma)} \le CM_d.$$

Proof. Consider a fixed $u_{ad} \in U_{ad}$ such that $u_{ad} \in U_{h,ad}$ for all h > 0. Using that $\|\bar{y}_h - y_d\|_{L^2(\Omega)}^2 \ge 0$ and the optimality of \bar{u}_h together with Young's

inequality and estimate (4.12), we have for all $\varepsilon > 0$ that

$$\begin{split} \frac{\nu}{2} \|\bar{u}_{h}\|_{L^{2}(\Gamma)}^{2} &\leq J_{h}(\bar{u}_{h}) - \int_{\Gamma} y_{h}(\bar{u}_{h})g_{\varphi} \,\mathrm{d}x \\ &\leq J_{h}(u_{\mathrm{ad}}) + \varepsilon \|y_{h}(\bar{u}_{h})\|_{L^{2}(\Gamma)}^{2} + \frac{1}{4\varepsilon} \|g_{\varphi}\|_{L^{2}(\Gamma)}^{2} \\ &\leq \frac{1}{2} \|y_{h}(u_{\mathrm{ad}}) - y_{d}\|_{L^{2}(\Omega)}^{2} + \frac{\nu}{2} \|u_{\mathrm{ad}}\|_{L^{2}(\Gamma)}^{2} + \int_{\Gamma} y_{h}(u_{\mathrm{ad}})g_{\varphi} \,\mathrm{d}x \\ &\quad + \varepsilon c_{2}^{2} \|\bar{u}_{h}\|_{L^{2}(\Gamma)}^{2} + \frac{1}{4\varepsilon} \|g_{\varphi}\|_{L^{2}(\Gamma)}^{2} \\ &\leq \|y_{h}(u_{\mathrm{ad}})\|_{L^{2}(\Omega)}^{2} + \|y_{d}\|_{L^{2}(\Omega)}^{2} + (\frac{\nu}{2} + c_{2}^{2})\|u_{\mathrm{ad}}\|_{L^{2}(\Gamma)}^{2} \\ &\quad + \varepsilon c_{2}^{2} \|\bar{u}_{h}\|_{L^{2}(\Gamma)}^{2} + \frac{1+\varepsilon}{4\varepsilon} \|g_{\varphi}\|_{L^{2}(\Gamma)}^{2} \end{split}$$

where c_2 is introduced in (4.12). Taking $\varepsilon = \nu/(4c_2^2)$, we readily deduce that $\{\bar{u}_h\}$ is uniformly bounded in $L^2(\Gamma)$. The estimate for $\|\bar{y}_h\|_{H^1(\Omega)}$ follows from this one and estimate (4.2).

Estimates for $\|\bar{u}\|_{L^2(\Gamma)}$ and $\|\bar{y}\|_{H^1(\Omega)}$ follow in a similar way. From this last one and Lemma 2.6 an estimate for $\|\bar{\varphi}\|_{H^1(\Omega)}$ in terms of the data is obtained. The trace theorem and the projection formula (5.1) lead to the estimate for $\|\bar{u}\|_{H^{1/2}(\Gamma)}$.

If Assumption 5.2 holds, then, using the estimate for $\|\bar{y}\|_{L^2(\Omega)}$ and noting that the condition $\beta_j < 1/2$ implies $\prod_{j=1}^m W_{\vec{\beta}}^{1/2,2}(\Gamma_j) \hookrightarrow L^2(\Gamma)$ and hence

 $\|y_d\|_{L^2(\Omega)} + \|g_{\varphi}\|_{L^2(\Gamma)} + 1 \le M_d,$

we obtain an estimate of $\|\bar{\varphi}\|_{W^{2,2}_{\bar{\beta}}(\Omega)}$ in terms of M_d . The trace theorem and the projection formula (5.1) lead to the estimate for $\|\bar{u}\|_{H^1(\Gamma)}$.

In the rest of the work s represents any positive number satisfying $s \leq 1$ and $s < \lambda_j / \mu_j$.

Lemma 5.6. Suppose Assumption 5.2 holds. Then, there exists C > 0 independent of h, y_d , g_{φ} and $\{\bar{u}_h\}$ such that

$$\|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{L^2(\Omega)} \le Ch^{3s/2} M_d.$$
 (5.3)

Moreover, for all $\theta \in (1/2, 1]$ we have the following estimate:

$$\|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{L^2(\Gamma)} \le Ch^{(2-\theta)s} M_d.$$
(5.4)

Proof. By the triangle inequality

$$\|\varphi_{\bar{u}_{h}} - \bar{\varphi}_{h}\|_{L^{2}(\Omega)} \le \|\varphi_{\bar{u}_{h}} - \varphi^{h}\|_{L^{2}(\Omega)} + \|\varphi^{h} - \bar{\varphi}_{h}\|_{L^{2}(\Omega)},$$
(5.5)

where φ^h is the unique element in $H^1(\Omega)$ such that $a(z, \varphi^h) = \int_{\Omega} (\bar{y}_h - y_d) z \, dx + \int_{\Gamma} g_{\varphi} z \, dx$ for all $z \in H^1(\Omega)$, i.e., $\bar{\varphi}_h$ is the finite element approximation of φ^h .

Let us estimate the first term in the right hand side of (5.5). Noting that

$$a(z,\varphi_{\bar{u}_h}-\varphi^h) = \int_{\Omega} (y_{\bar{u}_h}-y_h(\bar{u}_h)) z \, \mathrm{d}x \,\,\forall z \in H^1(\Omega),$$

we deduce from Theorem 3.5, the existence of C > 0 independent of h such that

$$\|\varphi_{\bar{u}_h} - \varphi^h\|_{L^2(\Omega)} \le C \|y_{\bar{u}_h} - y_h(\bar{u}_h)\|_{L^2(\Omega)}.$$
(5.6)

Applying the finite element error estimate for the state (4.19) of Corollary 4.6 and Lemma 5.5, we have

$$\begin{aligned} \|y_{\bar{u}_h} - y_h(\bar{u}_h)\|_{L^2(\Omega)} &\leq Ch^{3s/2} \|\bar{u}_h\|_{L^2(\Gamma)} \\ &\leq Ch^{3s/2} (\|y_d\|_{L^2(\Omega)} + \|g_{\varphi}\|_{L^2(\Gamma)} + 1) \leq Ch^{3s/2} M_d. \end{aligned}$$

This, together with (5.6) leads to

$$\|\varphi_{\bar{u}_h} - \varphi^h\|_{L^2(\Omega)} \le Ch^{3s/2} M_d.$$

$$(5.7)$$

To estimate the second summand in the right hand side of (5.5) we apply the finite element error estimate (4.15), the uniform boundness result in Lemma 5.5 and the embedding $\prod_{j=1}^{m} W_{\vec{\beta}}^{1/2,2}(\Gamma) \hookrightarrow L^2(\Gamma)$:

$$\begin{split} \|\varphi^{h} - \bar{\varphi}_{h}\|_{L^{2}(\Omega)} &\leq C \left(\|\bar{y}_{h} - y_{d}\|_{L^{2}(\Omega)} + \sum_{j=1}^{m} \|g_{\varphi}\|_{W^{1/2,2}_{\bar{\beta}}(\Gamma_{j})} \right) h^{2s} \\ &\leq C \left(\|\bar{y}_{h}\|_{L^{2}(\Omega)} + \|y_{d}\|_{L^{2}(\Omega)} + \sum_{j=1}^{m} \|g_{\varphi}\|_{W^{1/2,2}_{\bar{\beta}}(\Gamma_{j})} \right) h^{2s} \\ &\leq C \left(2\|y_{d}\|_{L^{2}(\Omega)} + \|g_{\varphi}\|_{L^{2}(\Gamma)} + \sum_{j=1}^{m} \|g_{\varphi}\|_{W^{1/2,2}_{\bar{\beta}}(\Gamma_{j})} + 1 \right) h^{2s} \\ &\leq C \left(\|y_{d}\|_{L^{2}(\Omega)} + \sum_{j=1}^{m} \|g_{\varphi}\|_{W^{1/2,2}_{\bar{\beta}}(\Gamma_{j})} + 1 \right) h^{2s} \\ &\leq C \left(\|y_{d}\|_{L^{2}(\Omega)} + \sum_{j=1}^{m} \|g_{\varphi}\|_{W^{1/2,2}_{\bar{\beta}}(\Gamma_{j})} + 1 \right) h^{2s} = Ch^{2s} M_{d}. \end{split}$$

Estimate (5.3) follows, hence, from (5.5) together with this last estimate and (5.7).

Let us prove (5.4). First we notice that for $1/2 < \theta \leq 1$, the trace operator is continuous from $H^{\theta}(\Omega)$ to $L^{2}(\Gamma)$, so

$$\|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{L^2(\Gamma)} \le C \|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{H^\theta(\Omega)}.$$

To estimate the term $\|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{H^{\theta}(\Omega)}$, we first introduce $\phi_h \in Y_h$, the finite element approximation of $\varphi_{\bar{u}_h}$, that satisfies $a(z_h, \phi_h) = \int_{\Omega} (y_{\bar{u}_h} - y_d) z_h \, dx + \int_{\Gamma} g_{\varphi} z_h \, dx$ for all $z_h \in Y_h$. The difference $\phi_h - \bar{\varphi}_h$ satisfies $a(z_h, \phi_h - \bar{\varphi}_h) = \int_{\Omega} (y_{\bar{u}_h} - y_h) z_h \, dx$

 $\int_{\Omega} (y_{\bar{u}_h} - \bar{y}_h) z_h \, dx$ for all $z_h \in Y_h$. From the continuity estimate for the discrete adjoint equation of Theorem 4.3 we deduce that

$$\|\phi_h - \bar{\varphi}_h\|_{H^1(\Omega)} \le C \|y_{\bar{u}_h} - \bar{y}_h\|_{L^2(\Omega)}.$$
(5.8)

Using the triangle inequality, the fact that $\theta \leq 1$, the finite element error estimate for the adjoint estate Eq. (4.19) of Corollary 4.6, (5.8), and the finite element error estimate for the state equation (4.20), together with the uniform boundness of $\|\bar{u}_h\|_{L^2(\Gamma)}$ provided in Lemma 5.5, we obtain

$$\begin{split} \|\varphi_{\bar{u}_{h}} - \bar{\varphi}_{h}\|_{H^{\theta}(\Omega)} &\leq \|\varphi_{\bar{u}_{h}} - \phi_{h}\|_{H^{\theta}(\Omega)} + \|\phi_{h} - \bar{\varphi}_{h}\|_{H^{1}(\Omega)} \\ &\leq C \left(h^{(2-\theta)s} M_{d} + \|y_{\bar{u}_{h}} - \bar{y}_{h}\|_{L^{2}(\Omega)} \right) \\ &\leq C \left(h^{(2-\theta)s} M_{d} + h^{3s/2} \|\bar{u}_{h}\|_{L^{2}(\Gamma)} \right) \leq C h^{(2-\theta)s} M_{d}, \end{split}$$

where the last inequality is a result of Lemma 5.5 and the condition $\theta > 1/2$.

We are now in position to prove the main result of this section.

Theorem 5.7. Suppose Assumption 5.2 holds. Then, there exists a constant independent of h, y_d and g_{φ} such that, for all $0 < h < \overline{h}$

$$\|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)} \le Ch^{s^*} M_d,$$

for all $s^* \leq 1$ such that $s^* < \frac{3}{2} \frac{\lambda_j}{\mu_j}$ for all $j \in \{1, \ldots, m\}$.

Proof. Testing the equality $a(z, \bar{\varphi} - \varphi_{\bar{u}_h}) = \int_{\Omega} (\bar{y} - y_{\bar{u}_h}) z \, dx$ for $z = \bar{y} - y_{\bar{u}_h}$ and using the state equation, we have that

$$0 \le \|\bar{y} - y_{\bar{u}_h}\|_{L^2(\Omega)}^2 = a(\bar{y} - y_{\bar{u}_h}, \bar{\varphi} - \varphi_{\bar{u}_h}) = \int_{\Gamma} (\bar{u} - \bar{u}_h)(\bar{\varphi} - \varphi_{\bar{u}_h}) \,\mathrm{d}x.$$

So we can write

$$\begin{split} \nu \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)}^2 &\leq \int_{\Gamma} (\bar{\varphi} - \varphi_{\bar{u}_h} + \nu(\bar{u} - \bar{u}_h))(\bar{u} - \bar{u}_h) \,\mathrm{d}x \\ &= \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_h + \nu(\bar{u} - \bar{u}_h))(\bar{u} - \bar{u}_h) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi}_h - \varphi_{\bar{u}_h})(\bar{u} - \bar{u}_h) \,\mathrm{d}x = I + II. \end{split}$$

Let us bound the first term. First we insert in appropriate places $Q_h \bar{u}$ and \bar{u} . Next, we apply the first order optimality conditions for the continuous and discrete problem. Finally we insert $\varphi_{\bar{u}_h}$ to obtain

$$\begin{split} I &= \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_{h} + \nu(\bar{u} - \bar{u}_{h}))(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_{h} + \nu(\bar{u} - \bar{u}_{h}))(Q_{h}\bar{u} - \bar{u}_{h}) \,\mathrm{d}x \\ &= \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_{h} + \nu(\bar{u} - \bar{u}_{h}))(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi} + \nu\bar{u})(Q_{h}\bar{u} - \bar{u}_{h}) \,\mathrm{d}x \\ &+ \int_{\Gamma} (\bar{\varphi}_{h} + \nu\bar{u}_{h}))(\bar{u}_{h} - Q_{h}\bar{u}) \,\mathrm{d}x \\ &= \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_{h} + \nu(\bar{u} - \bar{u}_{h}))(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi} + \nu\bar{u})(Q_{h}\bar{u} - \bar{u}) \,\mathrm{d}x \\ &+ \int_{\Gamma} (\bar{\varphi} + \nu\bar{u})(\bar{u} - \bar{u}_{h}) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi}_{h} + \nu\bar{u}_{h}))(\bar{u}_{h} - Q_{h}\bar{u}) \,\mathrm{d}x \\ &\leq \int_{\Gamma} (\bar{\varphi} - \bar{\varphi}_{h} + \nu(\bar{u} - \bar{u}_{h}))(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x + \int_{\Gamma} (\bar{\varphi} + \nu\bar{u})(Q_{h}\bar{u} - \bar{u}) \,\mathrm{d}x \\ &= \int_{\Gamma} (\bar{\varphi} - \varphi_{\bar{u}_{h}} + \nu(\bar{u} - \bar{u}_{h}))(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x + \int_{\Gamma} (\varphi_{\bar{u}_{h}} - \bar{\varphi}_{h})(\bar{u} - Q_{h}\bar{u}) \,\mathrm{d}x \\ &+ \int_{\Gamma} (\bar{\varphi} + \nu\bar{u})(Q_{h}\bar{u} - \bar{u}) \,\mathrm{d}x = I_{A} + I_{B} + I_{C}. \end{split}$$

From Lemmas 5.3 and 5.5, it is clear that $I_C \leq Ch^2 M_d^2$.

Let us study I_A . Testing the equality $a(z, \bar{\varphi} - \varphi_{\bar{u}_h}) = \int_{\Omega} (\bar{y} - y_{\bar{u}_h}) z \, dx$ for $z = \bar{y} - y_{Q_h \bar{u}}$ and using the state equation, Cauchy–Schwarz inequality, and Theorem 3.4(a), we obtain

$$\int_{\Gamma} (\bar{\varphi} - \varphi_{\bar{u}_h})(\bar{u} - Q_h \bar{u}) \, \mathrm{d}x = a(\bar{y} - y_{Q_h \bar{u}}, \bar{\varphi} - \varphi_{\bar{u}_h}) = \int_{\Omega} (\bar{y} - y_{\bar{u}_h})(\bar{y} - y_{Q_h \bar{u}}) \, \mathrm{d}x$$
$$\leq \|y_{\bar{u} - Q_h \bar{u}}\|_{L^2(\Omega)} \|y_{\bar{u} - \bar{u}_h}\|_{L^2(\Omega)} \leq C \|\bar{u} - Q_h \bar{u}\|_{L^2(\Gamma)} \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)}$$

Using this and Lemmas 5.3 and 5.5, we obtain

$$I_A = \int_{\Gamma} (\bar{\varphi} - \varphi_{\bar{u}_h} + \nu(\bar{u} - \bar{u}_h))(\bar{u} - Q_h \bar{u}) \, \mathrm{d}x \le Ch \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)} M_d.$$

Next we bound I_B and II. By the Cauchy-Schwarz inequality, we have that, for every $v \in L^2(\Gamma)$,

$$\int_{\Gamma} (\varphi_{\bar{u}_h} - \bar{\varphi}_h) v \, \mathrm{d}x \le \|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{L^2(\Gamma)} \|v\|_{L^2(\Gamma)}.$$
(5.9)

Taking $v = \bar{u} - Q_h \bar{u}$ in (5.9) and using (5.4) and Lemmas 5.3 and 5.5, we conclude that

$$I_B \le \|\varphi_{\bar{u}_h} - \bar{\varphi}_h\|_{L^2(\Gamma)} \|\bar{u} - Q_h \bar{u}\|_{L^2(\Gamma)} \le Ch^{(2-\theta)s+1} M_d^2.$$

Finally, taking $v = \bar{u} - \bar{u}_h$ in (5.9) and using (5.4), we have

$$II \le Ch^{(2-\theta)s} \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)} M_d.$$

Gathering all the estimates we have that

$$\nu \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)}^2 \le C(h\|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)}M_d + h^{(2-\theta)s+1}M_d^2 + h^2M_d^2 + h^{(2-\theta)s}\|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)}M_d)$$

and the proof concludes using Young's inequality. Notice that the appearance of the terms $h \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)} M_d$ and $h^2 M_d^2$ implies that the resulting exponent s^* is less or equal than one. On the other hand, since $\theta > 1/2$, the term $h^{(2-\theta)s} \|\bar{u} - \bar{u}_h\|_{L^2(\Gamma)} M_d$ yields the bound $s^* \leq (2-\theta)s < \frac{3}{2}s < \frac{3}{2}\frac{\lambda_j}{\mu_j}$. Finally, from the term $h^{(2-\theta)s+1}M_d^2$ we obtain the bound $s^* \leq \min\{(2-\theta)s,1\}$, so no new conditions are imposed on s^* .

6. A Numerical Example

Let Ω be the *L*-shaped domain $\Omega = \{x \in \mathbb{R}^2 : r < \sqrt{2}, \theta < 3\pi/2\} \cap (-1, 1)^2$. We consider a functional of the form

$$J(u) = \frac{1}{2} \int_{\Omega} (y_u(x) - y_d(x))^2 \, \mathrm{d}x + \frac{\nu}{2} \int_{\Gamma} u(x)^2 \, \mathrm{d}x + \int_{\Gamma} y_u(x) g_{\varphi}(x) \, \mathrm{d}x,$$

where

$$\begin{cases} -\Delta y_u + b \cdot \nabla y_u + a_0 y_u = f & \text{in } \Omega, \\ \partial_n y = u + g_y & \text{on } \Gamma. \end{cases}$$

with data ν , y_d , g_{φ} , b, a_0 , g_y described below. The inclusion of data f and g_y is useful to write a problem with known exact solution. Notice that, if we denote $y_0 \in L^2(\Omega)$ the state related to $u \equiv 0$ and redefine $y_d := y_d - y_0$ and $y_u := y_u - y_0$, the problem fits into the framework of problem (P) and Eq. (1.1).

Let (r, θ) be the polar coordinates in the plane, $r \ge 0$, $\theta \in [0, 2\pi]$. The interior angle at the vertex of the domain located at the origin is $\omega = \omega_1 = 3\pi/2$ and we denote $\lambda = \lambda_1 = \pi/\omega_1 = 2/3$. For $j = 2, \ldots, 6$, $\omega_j = \pi/2$ and $\lambda_j = 2$.

We introduce $\bar{y} = r^{\lambda} \cos(\lambda \theta)$, $\bar{\varphi} = -\bar{y}$ and $\bar{u} = -\bar{\varphi}/\nu$ on Γ and, for some $\alpha > -3/2$ and some $\delta \ge 0$, we consider $b(x) = \delta r^{\alpha+1} (\cos \theta, \sin \theta)^T$ and $a_0(x) = r^{\alpha}$.

The data for this problem are defined as $f = b \cdot \nabla \bar{y} + a_0 \bar{y}$, $g_y = \partial_n \bar{y} - \bar{u}$ on Γ , $y_d = \bar{y} + \nabla \cdot (\bar{\varphi}b) - a_0 \bar{\varphi}$ and $g_{\varphi} = \partial_n \bar{\varphi} + (b\bar{\varphi}) \cdot n$.

For all $\alpha > -2$, $b \in L^{\hat{p}}(\Omega)$ for some $\hat{p} > 2$ (Assumption 2.1). For $\alpha > -1-\beta$, $a_0, \nabla \cdot b, f, y_d \in L^2_{\beta}(\Omega)$ and $b \cdot n, g_y, g_{\varphi} \in W^{1/2,2}_{\beta}(\Gamma)$, so the assumptions of Theorems 3.4(c) and 3.5(c) hold. If we impose $\beta < 1/2$ (assumption in Theorem 5.1), we have that for $\alpha > -3/2$ all the assumptions of the paper hold. In our experiments, we fix $\alpha = -1.25$.

The given \bar{u} is the solution of the control problem

(P)
$$\min_{u \in L^2(\Gamma)} J(u),$$

with related state \bar{y} and adjoint state $\bar{\varphi}$, which satisfy the optimality system

$$\begin{cases} -\Delta \bar{y} + b \cdot \nabla \bar{y} + a_0 \bar{y} = f & \text{in } \Omega, \\ \partial_{n_A} \bar{y} = g_y + \bar{u} & \text{on } \Gamma, \end{cases} \\ \begin{cases} -\Delta \bar{\varphi} - \nabla \cdot (b\bar{\varphi}) + a_0 \bar{\varphi} = \bar{y} - y_d & \text{in } \Omega, \\ \partial_n \bar{\varphi} + \bar{\varphi} b \cdot n = g_{\varphi} & \text{on } \Gamma, \end{cases} \\ \bar{u} = -\bar{\varphi}/\nu & \text{on } \Gamma. \end{cases}$$

It is clear that $\bar{y}, \bar{\varphi} \in W^{2,2}_{\vec{\beta}}(\Omega)$ and $\bar{u} \in H^1(\Gamma) \cap W^{1/2,2}_{\vec{\beta}}(\Gamma)$ for $\vec{\beta} = (\beta, 0, 0, 0, 0, 0, 0)$ for all $\beta > 1 - \lambda > 1/3$.

For $\delta = 6$, we have checked numerically that the operator is not coercive,

To discretize the problem we use the finite element approximation described in the work. We use a family of graded meshes obtained by bisection; see, e.g., [25, Figure 1.2]. This meshing method does not lead to superconvergence properties in the gradients. The code has been done with Matlab on a desktop PC with Interl(R) Core(TM) i5-7500CPU at 3.4GHz with 24GB of RAM. The meshes have been prepared using functions provided by Johannes Pfefferer. The finite element approximations are obtained with code prepared by us and the linear systems are solved using Matlab's [L,U,P,Q,D] = lu(S)method. The optimization of the resulting finite-dimensional quadratic program is done using Matlab's pcg.

First we check estimates (4.14) and (4.15) for the error in the solution of the boundary value problem. For appropriately graded meshes, $\mu < 2/3 = \lambda$, we expect order h^2 in $L^2(\Omega)$ and order h in $H^1(\Omega)$. For a quasi-uniform family, $\mu = 1$, we have s < 2/3, so we expect order $h^{1.33}$ in $L^2(\Omega)$ and order $h^{0.66}$ in $H^1(\Omega)$. We summarize the results in Tables 1, 2, 3 and 4. We include results for both the state and adjoint state equation. Notice that $\tilde{\varphi}_h$ is the finite element approximation of $\bar{\varphi}$, obtained using the exact \bar{y} , i.e., $a(z_h, \tilde{\varphi}_h) = \int_{\Omega} (\bar{y} - y_d) z_h \, dx + \int_{\Gamma} g_{\varphi} z_h \, dx$ for all $z_h \in Y_h$.

Next, we turn to the control problem and check the estimate in Theorem 5.7. Notice that we should obtain order of convergence h for both graded-meshes and quasi-uniform meshes. We summarize the results in Table 5.

Note that in this example the regularity of the adjoint state is even $\bar{\varphi} \in W^{2,\infty}_{\vec{\gamma}}(\Gamma)$ for $\vec{\gamma} = (\gamma, 0, 0, 0, 0, 0)$ with $\gamma > 4/3$. This leads to superconvergence properties in the convergence in the norms of $L^2(\Omega)$ and $L^2(\Gamma)$ of both the state and adjoint state variable, where, despite expecting order of convergence 1, as for the control, we obtain the same order of convergence as the one for the boundary value problem, i.e., 1.33 or almost 2 in our examples. This phenomenon will be studied in a future paper.

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2.92e-01 1.94e-01 1.95e-01		$\ \varphi - \varphi_h\ _{L^2(\Omega)}$	EOC	$\ \bar{\varphi} - \tilde{\varphi}_h\ _{H^1(\Omega)}$	EOC
1.94e-01 1.95 $a-01$		$3.85e{-}02$		$3.28e{-01}$	
1 950-01	0.59	$1.27e{-}02$	1.60	2.02e-01	0.70
T 00-21).63	$4.45e{-}03$	1.51	$1.27e{-01}$	0.68
7.98e - 02	0.65	1.64e - 03	1.44	7.98e - 02	0.67
$5.05e{-}02$).66	6.24e - 04	1.39	5.04e - 02	0.66
$3.19\mathrm{e}{-02}$).66	$2.42\mathrm{e}{-04}$	1.37	$3.19\mathrm{e}{-02}$	0.66
$2.02\mathrm{e}{-02}$).66	$9.47\mathrm{e}{-05}$	1.35	2.01e-02	0.66
$1.27e{-}02$).66	$3.73\mathrm{e}{-05}$	1.34	$1.27e{-}02$	0.66
$8.02\mathrm{e}{-03}$	0.67	$1.47\mathrm{e}{-05}$	1.34	$8.01e{-}03$	0.66
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8.02e-03 le, orders of convergence ex).67).66 pected	면	1.47e-05 from (4.14). Quasi-	1.47e-05 1.34 1.33 from (4.14). Quasi-uniform	$\begin{array}{cccccccccccccccccccccccccccccccccccc$

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		$\ \varphi - \varphi_h\ _{L^2(\Omega)}$	EOO	$\ \varphi - \varphi h\ _{H^1(\Omega)}$	EOC
$2.48\mathrm{e}{-01}$		$2.27\mathrm{e}{-02}$		$2.65e{-01}$	
$1.40e{-01}$	0.82	$5.53\mathrm{e}{-03}$	2.03	$1.42\mathrm{e}{-01}$	0.90
$7.74e{-}02$	0.86	$1.50\mathrm{e}{-03}$	1.89	$7.76\mathrm{e}{-02}$	0.87
$4.22e{-}02$	0.88	4.20e - 04	1.83	$4.22e{-}02$	0.88
$2.27\mathrm{e}{-02}$	0.89	1.18e - 04	1.83	$2.27\mathrm{e}{-02}$	0.89
$1.21e{-}02$	0.91	$3.28\mathrm{e}{-05}$	1.84	1.21e-02	0.91
$6.42e{-03}$	0.92	$9.07\mathrm{e}{-06}$	1.85	$6.42e{-}03$	0.92
$3.38e{-03}$	0.92	$2.49\mathrm{e}{-06}$	1.86	$3.38e{-03}$	0.92
$1.77\mathrm{e}{-03}$	0.93	6.80e - 07	1.87	$1.77e{-}03$	0.93
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	7.74e-02 4.22e-02 2.27e-02 1.21e-02 6.42e-03 3.38e-03 1.77e-03	$\begin{array}{cccccc} 7.74\mathrm{e}{-}02 & 0.86 \\ 4.22\mathrm{e}{-}02 & 0.88 \\ 2.27\mathrm{e}{-}02 & 0.89 \\ 1.21\mathrm{e}{-}02 & 0.91 \\ 6.42\mathrm{e}{-}03 & 0.92 \\ 3.38\mathrm{e}{-}03 & 0.92 \\ 1.77\mathrm{e}{-}03 & 0.93 \\ 1.77\mathrm{e}{-}03 & 1.93 \end{array}$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$

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			FCC	$\ \varphi - \varphi h\ _{H^1(\Omega)}$	EUC
8e-01		$2.27e{-}02$		$2.65e{-01}$	
Je−01	0.82	$5.53\mathrm{e}{-03}$	2.03	$1.42e{-}01$	0.90
4e-02	0.99	1.28e - 03	2.12	$7.05e{-}02$	1.01
4e-02	0.91	3.46e - 04	1.88	$3.74e{-}02$	0.91
)e-02	0.90	$9.50\mathrm{e}{-05}$	1.86	2.00e-02	0.90
1e-02	0.99	$2.42\mathrm{e}{-05}$	1.97	1.01e-02	0.99
2e-03	0.95	$6.34e{-}06$	1.93	$5.22\mathrm{e}{-03}$	0.95
1e-03	0.95	$1.67\mathrm{e}{-06}$	1.92	$2.71\mathrm{e}{-03}$	0.95
3e-03	0.99	$4.22\mathrm{e}{-07}$	1.99	$1.36e{-}03$	0.99
			2		1
4e02 De02 Le03 Le03 Se03 Se03	$\begin{array}{c} 0.91 \\ 0.90 \\ 0.95 \\ 0.95 \\ 0.95 \end{array}$	9 - 1 - 0 - 7 - 7 - 7	46e-04 50e-05 42e-05 34e-06 67e-06 22e-07	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{rrrrr} 46e-04 & 1.88 & 3.74e-02 \\ 50e-05 & 1.86 & 2.00e-02 \\ 42e-05 & 1.97 & 1.01e-02 \\ 34e-06 & 1.93 & 5.22e-03 \\ 67e-06 & 1.92 & 2.71e-03 \\ 22e-07 & 1.99 & 1.36e-03 \end{array}$

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	$\ \bar{y} - \bar{y}_h(\bar{u})\ _{L^2(\Omega)}$	EOC	$\ \bar{y} - \bar{y}_h(\bar{u})\ _{H^1(\Omega)}$	EOC	$\ \bar{\varphi} - \tilde{\varphi}_h\ _{L^2(\Omega)}$	EOC	$\ ar{arphi}- \widetilde{arphi}_h\ _{H^1(\Omega)}$	EOC
	$8.30e{-}02$		$2.48e{-01}$		$2.27\mathrm{e}{-02}$		$2.65\mathrm{e}{-01}$	
7	$1.80e{-}02$	2.20	$1.26e{-01}$	0.98	$4.50\mathrm{e}{-03}$	2.33	$1.27e{-}01$	1.06
3	$4.20 \mathrm{e}{-03}$	2.10	$6.31e{-}02$	0.99	$1.10e{-}03$	2.03	$6.32e{-}02$	1.00
4	$1.07\mathrm{e}{-03}$	1.97	$3.19e{-}02$	0.99	$2.83e{-}04$	1.96	$3.19e{-}02$	0.99
5	$2.75e{-}04$	1.96	$1.61e{-}02$	0.99	$7.22\mathrm{e}{-05}$	1.97	$1.61e{-}02$	0.99
9	$7.03\mathrm{e}{-05}$	1.97	$8.11e{-}03$	0.99	$1.83\mathrm{e}{-05}$	1.98	$8.11e{-}03$	0.99
2	$1.78\mathrm{e}{-05}$	1.98	4.08e - 0.3	0.99	$4.60\mathrm{e}{-06}$	1.99	4.08e - 03	0.99
x	$4.51\mathrm{e}{-06}$	1.98	$2.05e{-}03$	0.99	1.16e - 06	1.99	$2.05\mathrm{e}{-03}$	0.99
6	$1.14e{-}06$	1.99	$1.03e{-}03$	1.00	$2.90\mathrm{e}{-07}$	2.00	$1.03\mathrm{e}{-03}$	1.00
Exp	ected	2		1		2		1
	In the I	ottom line	e, orders of convergen	ce expecte	ed from (4.14). Gre	aded mesh	$\mu = 0.5$	
				I	~			

ders of convergence for the optimal control problem
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TABLE 5.

.—	Quasi-uniform $\ \bar{u} - \bar{u}_h\ _{L^2(\Gamma)}$	EOC	$\mu = 0.66$ $\ \bar{u} - \bar{u}_{h}\ _{L^{2}(\Gamma)}$	EOC	$\mu = 0.6$ $\ \bar{u} - \bar{u}_{h}\ _{L^{2}(\Gamma)}$	EOC	$\mu = 0.5$ $\ \bar{u} - \bar{u}_h\ _{L^2(\Gamma)}$	EOC
, <u> </u>	2.58e-01		2.58e-01		2.58e-01		2.58e-01	
2	$1.33e{-01}$	0.96	1.15e - 01	1.16	$1.15e{-01}$	1.16	$1.15e{-01}$	1.17
°	$6.79\mathrm{e}{-02}$	0.97	6.05e - 02	0.93	$5.62\mathrm{e}{-02}$	1.03	$5.61\mathrm{e}{-02}$	1.03
4	$3.45\mathrm{e}{-02}$	0.98	2.90e - 02	1.06	$2.90\mathrm{e}{-02}$	0.96	$2.80\mathrm{e}{-02}$	1.01
r.	$1.75\mathrm{e}{-02}$	0.98	1.48e - 02	0.97	$1.44e{-}02$	1.00	$1.41e{-}02$	0.98
9	$8.80e{-03}$	0.99	$7.34e{-}03$	1.01	$7.22\mathrm{e}{-03}$	1.00	$7.07e{-}03$	1.00
2	$4.42e{-}03$	0.99	3.68e - 03	1.00	$3.61\mathrm{e}{-03}$	1.00	$3.53\mathrm{e}{-03}$	1.00
∞	$2.22e{-}03$	1.00	1.84e - 03	1.00	$1.80e{-}03$	1.00	$1.76e{-}03$	1.00
6	$1.11e{-}03$	1.00	$9.19e{-}04$	1.00	$9.01\mathrm{e}{-04}$	1.00	$8.82e{-}04$	1.00
Exp	ected	1		1		1		1

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Data availability Not applicable.

Declarations

Conflict of interest No conflict of interest to disclose.

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Consent to participate Not applicable.

Consent for publication Not applicable.

Code availability Not applicable.

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Thomas Apel Institute of Mathematics and Computer-Based Simulation Universität der Bundeswehr München 85577 Neubiberg Germany e-mail: thomas.apel@unibw.de

Mariano Mateos Departamento de Matemáticas Universidad de Oviedo 33203 Gijón Spain e-mail: mmateos@uniovi.es

Arnd Rösch Fakultät für Mathematik Universtät Duisburg-Essen 45127 Essen Germany e-mail: arnd.roesch@uni-due.de

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